1 What is the skill of ocean tracers in reducing uncertainties

2 about ocean diapycnal mixing and projections of the Atlantic

3 Meridional Overturning Circulation?

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- 6 Received 26 May 2010; revised 30 July 2010; accepted 25 August 2010; published XX Month 2010.
- 7 [1] Current projections of the oceanic response to anthropogenic climate forcings are
- 8 uncertain. Two key sources of these uncertainties are (1) structural errors in current Earth
- 9 system models and (2) imperfect knowledge of model parameters. Ocean tracer
- 10 observations have the potential to reduce these uncertainties. Previous studies typically
- 11 consider each tracer separately, neglect potentially important statistical properties of the
- 12 system, or use methods that impose rather daunting computational demands. Here we
- 13 extend and improve upon a recently developed approach using horizontally averaged
- 14 vertical profiles of chlorofluorocarbon (CFC-11), radiocarbon (Δ^{14} C), and temperature (T)
- 15 observations to reduce model parametric and structural uncertainties. Our method
- 16 estimates a joint probability density function, which considers cross-tracer correlations and
- 17 spatial autocorrelations of the errors. We illustrate this method by estimating two model
- 18 parameters related to the vertical diffusivity, the background vertical diffusivity, and the
- 19 upper Southern Ocean mixing. We show that enhancing the upper Southern Ocean mixing in
- 20 the model improves the representations of ocean tracers and improves the hindcasts of the
- 21 Atlantic Meridional Overturning Circulation (AMOC). The most probable value of the
- 22 background vertical diffusivity in the pelagic pycnocline is between 0.1 and 0.2 cm² s⁻¹.
- 23 According to the statistical method, observations of Δ^{14} C reduce the uncertainty about the
- 24 background vertical diffusivity mostly followed by CFC-11 and T. Using all three tracers
- 25 jointly reduces the model uncertainty by 40%, more than each tracer individually. Given
- 26 several important caveats, we illustrate how the reduced model parametric uncertainty
- 27 improves probabilistic projections of the AMOC.
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31 1. Introduction

32 [2] The North Atlantic Overturning Circulation (AMOC) 33 is a key component of the climate system [Munk and 34 Wunsch, 1998]. Past changes in the AMOC intensity are 35 associated with considerable changes in global scale tem-36 perature and precipitation patterns [McManus et al., 2004]. 37 Anthropogenic climate forcings may trigger an AMOC 38 threshold response, with potentially serious impacts on

2007; *Keller et al.*, 2000]. Current AMOC model predic- 40 tions are deeply uncertain [*Zickfeld et al.*, 2007; *Meehl et al.*, 41 2007]. 42
[3] Tracer observations such as chlorofluorocarbon-11 43

natural systems and human welfare [Patwardhan et al., 39

- [3] Tracer observations such as chlorofluorocarbon-11 43 (CFC-11) and radiocarbon (Δ^{14} C) provide information on 44 the ventilation rate and advective properties in the ocean on 45 time scales ranging from decadal to centennial that can be 46 used for evaluating the skill of climate models in simulating 47 the ocean circulation [Doney et al., 2004]. A better representation of these processes in models can possibly improve 49 AMOC projections.
- [4] A key variable for determining ocean circulation 51 properties in models is the vertical ocean diffusivity (K_v) . 52 Changing this value in model simulations has a large impact 53 on oceanic heat storage and transport, uptake of ocean tra- 54 cers such as CO_2 [Sokolov et al., 1998], and on the work 55 necessary to lift the abyssal waters through stratification 56 (that closes the MOC circulation) [Wunsch and Ferrari, 57
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58 2004]. This variable is highly uncertain [Munk and Wunsch, 59 1998], and it is sometimes tuned in models to generate a 60 realistic AMOC strength [Gao et al., 2003]. In addition, this 61 parameter value affects the existence of multiple states of 62 the MOC in model simulations [Schmittner and Weaver,

[5] Various processes lead to mixing in the ocean such as 65 shear or buoyancy forced turbulence, interactions of flow 66 with topography, and double diffusion (differential molec-67 ular diffusion of heat and salt). See Smyth and Moum [2001] 68 and Moum and Smyth [2001] for reviews. Although General 69 Circulation Models have been increasing their ability of 70 parameterizing subgrid scale turbulent processes in the 71 ocean [Bryan and Lewis, 1979; Pacanowski and Philander, 72 1981; Large et al., 2004; Ferrari et al., 2008], due to the 73 complexity of the problem and processes involved, most 74 schemes are still highly simplified and parameterized. In 75 Earth System Models of Intermediate Complexity (EMICs), 76 the absence of more complex parameterizations elevates the 77 importance of the parameters related to K_{ν} in order to fulfill 78 the model necessity of turbulent mixing in simulating a 79 realistic AMOC strength.

[6] Several studies [e.g., England, 1993; Gao et al., 2003] 81 analyze the importance of the magnitude of the diffusivity 82 strength and parameterization on the MOC structure and 83 representations of tracers in ocean models. These studies are 84 typically silent on the question of how much information is 85 contained in the different types of observations. This is an 86 important question, for example, to inform the design of 87 AMOC observation and prediction systems [cf. Baehr et al., 88 2008; Keller et al., 2007].

[7] Schmittner et al. [2009] discusses a relatively simple 90 but computationally efficient method to estimate the back-91 ground ocean diffusivity K_{bg} from the combination of spa-92 tially resolved ocean tracer observations considering both, 93 observational and model errors. However, Schmittner et al. 94 [2009] neglects the effects of cross correlation between 95 different tracers, which limits the number of tracers that can 96 be combined in a joined probability density function. In 97 another recent study, Bhat et al. [2009] estimates the posterior 98 probability distribution for K_{bg} using Δ^{14} C and CFC-11 99 observations. Their approach uses a Gaussian process 100 emulator for the climate model and estimates the distribution 101 of K_{bg} via a Bayesian approach. While their kernel mixing 102 based approach to constructing the emulator is flexible and 103 efficient, it is conceptually complex and computationally 104 highly demanding for routine use with more than two ocean 105 tracers.

[8] Here we estimate the probability density function 107 (pdf) of K_{bg} using three tracers simultaneously. Our 108 approach provides a fast and easy way to implement the 109 methodology, enabling the routine use of information from 110 several ocean tracers jointly, while still considering spatial 111 autocorrelation as well as cross correlation between residuals 112 of different tracers. We demonstrate how neglecting cross 113 correlation and/or simplifying the mean function can com-114 promise the accuracy of the estimation. We improve the 115 treatment of uncertainties surrounding K_{ν} in the model by 116 considering the structural uncertainty about the upper 117 Southern Ocean mixing (u_K_{SO}) . We show that an ensemble 118 with enhanced Southern Ocean mixing is more consistent 119 with the observations.

[9] Furthermore, we advance on previous work by quan- 120 tifying and ranking the skill of the tracers CFC-11, Δ^{14} C 121 and temperature (T) to constrain the uncertainties in the 122 model parameter K_{bg} . We demonstrate the potential utility 123 of the considered observations to improve model predictions 124 of the AMOC.

2. Methods

Earth System Model of Intermediate Complexity

[10] We use the University of Victoria Earth System Model 128 of Intermediate Complexity (UVic 2.8) [Weaver et al., 2001]. 129 This model has been widely used in climate simulations and 130 models comparisons studies. In the UVic model, we parameterize the diapycnal diffusivity as $K_v = K_{tidal} + K_{SO} + K_{bg}$, which consists of the diffusivity due to local dissipation of 133 tidal energy and its resulting generation of turbulence and 134 mixing (K_{tidal}) [Simmons et al., 2004], a parameterization 135 for the vigorous mixing (K_{SO}) observed in the Southern 136 Ocean [e.g., Naveira Garabato et al., 2004], plus a back- 137 ground diffusivity K_{bg} that represents all other processes that 138 lead to mixing, such as nonlocal dissipation of tidal energy, 139 mesoscale eddy activity, double diffusion, hurricanes, 140 interaction of flow with topography, and others.

[11] The model accounts for increased mixing over rough 142 topography based on the tidal mixing scheme of St. Laurent 143 et al. [2002], and uses the [Gent and McWilliams, 1990] 144 eddy mixing parameterization. It is likely that K_{bg} is spa- 145tially and temporally variable in nature [Sriver et al., 2010], 146 but due to a lack of a more explicit representation of the 147 processes and for simplicity we assume a constant value of 148 K_{bg} everywhere. Note that K_{tidal} decays exponentially (with 149) an e-folding depth of 500 m above the seafloor) such that it 150 is unimportant in the pelagic pycnocline (i.e., away from the 151 boundaries). However, it is the value of K_{bg} in the pelagic 152 pycnocline that is most important in determining the large- 153 scale ocean circulation in models [cf., Marotzke, 1997; Munk 154 and Wunsch, 1998]. For the Southern Ocean (south of 40S) 155 parameterization, the vertical mixing is truncated at 1 cm²/s as 156 a lower bound ($K_v > 1 \text{ cm}^2/\text{s}$). The Southern Ocean is one of 157 the most tempestuous oceans on Earth, and these transient 158 effects may produce strong turbulent mixing, specially in the 159 upper Southern Ocean. In order to include uncertainties 160 about the upper Southern Ocean mixing, we further divide 161 the Southern Ocean mixing into upper (u_K_{SO}) and lower 162 (l_K_{SO}) parts. Therefore, $K_{SO} = u_K_{SO} + l_K_{SO}$, where 163 u_K_{SO} is the Southern Ocean mixing in the upper 500 m, 164 and l K_{SO} is the Southern Ocean mixing from 500m to the 165 bottom of the water column.

[12] We create two ensembles to analyze the uncertainty 167 in two model parameters, the background ocean diffusivity 168 (K_{bg}) and the upper Southern Ocean diffusivity (u K_{SO}). 169 Each ensemble contains seven members, corresponding to a 170 grid of the parameter K_{bg} values of (0.05, 0.1, 0.15, 0.2, 0.3, 171 0.4, and 0.5) cm² s⁻¹. The difference between the two 172 ensembles is that in the first one (ENSEMBLE 1), the 173 enhanced SO mixing is only applied in the lower part of the 174 Southern Ocean, so in the upper SO the mixing is equal to the 175 rest of the pelagic areas of the upper ocean (with indices 176 $u_{K_{SO}} = 0$, $l_{K_{SO}} = 1$), whereas the second one (ENSEMBLE 177 2) uses an enhanced mixing in the entire column of the 178 Southern Ocean (with indices u $K_{SO} = 1$, l $K_{SO} = 1$). As we 179

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180 are not varying the parameter $l_{-}K_{SO}$, it is suppressed in the 181 rest of the manuscript.

[13] The ocean component in UVic is MOM2 183 [Pacanowski, 1995] with a $1.8^{\circ} \times 3.6^{\circ}$ resolution in the 184 horizontal and 19 depth levels. The atmospheric component 185 is a one-layer atmospheric energy-moisture balance model, 186 which does not apply flux correction and is forced by pre-187 scribed winds from the NCAR/NCEP climatology. Also 188 included in the model are a thermodynamic sea ice com-189 ponent, a terrestrial vegetation (TRIFFID), and an oceanic 190 biogeochemistry based on the ecosystem model of 191 Schmittner et al. [2005].

[14] A total of 47,600 model years was preformed, what 193 makes UVic suitable for this kind of study. At first, the 194 model is spun up from observed data fields as initial con-195 ditions for 3000 years (with a coupled carbon cycle for the 196 last 1000 years) for each parameter value. It is then inte-197 grated from years 1800–2100 using historical and projected 198 climate forcings (SRES-A1FI scenario), extended to the 199 year 2200 following [Zickfeld et al., 2008]. We modify the 200 model to include non-CO₂ greenhouse gases, volcanic and 201 sulfate forcings from Sato et al. [1993] and Hansen and 202 Sato [2004]. Atmospheric sulfates data enter the model as 203 gridded optical depth [Koch et al., 1999], and follow the 204 same rate of decrease as the CO₂ concentration after 2100.

205 **2.2. Data**

[15] We focus on a subset of observations that have pre-207 viously been shown to provide constraints on the parame-208 terization of K_{ν} in ocean models: (1) temperature (T), (2) 209 chlorofluorocarbon 11 (CFC-11), and (3) radiocarbon 210 (Δ^{14} C) observations [cf. Schmittner et al., 2009; Bhat et al., 211 2009; Toggweiler et al., 1989]. Δ^{14} C is defined as the 212 ¹⁴C/¹²C ratio of air-sea fractionation-corrected data [Stuiver 213 and Polach, 1977]. Each of the tracers in this subset has a 214 different behavior and can constrain K_{ν} in different ways. 215 The temperature observations constrain K_{ν} , because K_{ν} 216 affects, for example, the shape of the thermocline as well as 217 the penetration of the anthropogenic heat anomalies 218 [Gnanadesikan, 1999]. The Δ^{14} C observations can con-219 strain K_{ν} in two main ways, because it has a natural and an 220 anthropogenic component. The natural component can 221 provide information of mixing rates (that are, in turn, a 222 function of K_{ν}) in the order of centuries or millennia. The 223 anthropogenic component, which greatly increased during 224 the 1950s and 1960s due to thermonuclear explosions, 225 provides information on decadal time scale. Here we do not 226 make distinction between natural and bomb ¹⁴C, thus we use 227 its total concentration. The anthropogenic tracer CFC-11 228 also constrains K_{ν} on decadal time scale, because atmo-229 spheric emissions started in the 1930s. The solubility of 230 CFCs in water is dependent on the temperature. Considering 231 CFC-11 and Δ^{14} C jointly can provide new insights into 232 vertical oceanic mixing because they have very different 233 forcing histories, air-sea equilibration time scales and water 234 solubility [Broecker and Peng, 1987; Ito et al., 2004], and 235 the observation errors and signal-to-noise ratios of the two 236 tracers are different. We analyze published data products for 237 these three tracers [Locarnini et al., 2006; Key et al., 2004] 238 and average the model hindcasts over the time the observa-239 tions have been collected, i.e., 1990s for CFC-11 and Δ^{14} C, 240 and 1950-2000 for temperature. We interpolate the observations to the model grid and the model output is restricted to 241 the regions where the data products are available. All con- 242 sidered ocean tracer observations are horizontally averaged 243 into global mean vertical profiles. Further, the probability 244 distributions of the model parameters, inferred from the 245 information of ocean tracers profiles, are compared with the 246 distribution inferred from the climatological observations of 247 the AMOC strength at 24°N. For this purpose, we use the 248 information of the AMOC strength calculated with the 249 inverse model of Lumpkin and Speer [2003], which is esti- 250 mated as $(17.6 \pm 2.7 \text{ Sv})$. The model ensembles are cali- 251 brated against observations using a Bayesian inference 252 method. We assume a Gaussian likelihood function and 253 estimate the posterior probability of K_{bg} and $u_{\underline{K}SO}$ given 254 the observations through a Markov Chain Monte Carlo 255 (MCMC) method [Metropolis et al., 1953]. Our method 256 accounts for autocorrelations of the residuals, as well as 257 cross correlation between residuals of different tracers. For 258 this, a separable covariance matrix Σ is estimated. The 259 inversion and the numerical implementation of the calibration 260 procedure are detailed in the next subsection. Readers not 261 interested in the details of the statistical inversion technique 262 can skip the next subsection without loss of understanding.

2.3. Bayesian Model Inversion

[16] The goal of Bayesian parameter estimation is to infer 265 a probability distribution(s) $p(\theta \mid O)$ representing the uncer- 266 tainty in one (or more) climate model parameter θ , condi- 267 tional on a vector of observed data O. Here θ are parameters 268 K_{be} and u K_{SO} , which are related to the vertical ocean dif- 269 fusivity in UVic. The inferential procedure is based on a 270 statistical model that relates the model parameters (θ) to the 271 observations (O) by way of the ensemble of model output M 272 (θ) . The statistical model used here assumes that the ob- 273 servations are randomly distributed around the model pre- 274 diction, according to

$$O = M(\theta) + \epsilon, \tag{1}$$

where the error is a random variable drawn from a multi- 276 variate normal distribution 277

$$\epsilon \sim N(\mu, \Sigma),$$
 (2)

with an unknown mean or bias term μ and covariance matrix 278 Σ . These distributional parameters are estimated along with 279 the model parameter θ . The error term encompasses all 280 processes which may cause the observations to deviate from 281 the model predictions, including model structural error, 282 unresolved variability in the climate system, and measure- 283 ment error. We model these errors as random processes, 284 approximated here by a potentially correlated Gaussian 285 probability function.

[17] The error mean term μ represents model bias, which 287 is common for each observed variable across ensemble 288 members. Schmittner et al. [2009] assumed a bias which is 289 constant with depth. Here we expand upon this form by using a general linear form that varies with depth (z), $\mu = az + 291$ b. This form improves the model fit as indicated by a 292 exploratory data analysis in the next section. The covariance 293 matrix, described later, captures the residual variability that is 294 unaccounted by the linear bias term.

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[18] The above probability model describing the spread of 297 observations about the model output defines a likelihood 298 function $L(O|\theta, \mu, \Sigma)$ for the data conditional on the model 299 and covariance parameters

$$L(O|\theta, \mu, \Sigma) = (2\pi)^{-N/2} |\Sigma|^{-1/2} \exp\left(-\frac{1}{2}\tilde{r}^T \Sigma^{-1}\tilde{r}\right), \quad (3)$$

300 where Σ is a covariance matrix and $\tilde{r} = O - M(\theta) - \mu$ are the 301 bias-corrected data-model residuals.

[19] Consider an ensemble M containing p runs of a cli-303 mate model, where each run corresponds to a different value 304 of a climate model parameter, θ_k , k = 1, ..., p. For each 305 ensemble member we analyze n ocean tracer profiles 306 defined at d spatial locations (depths). The matrix Σ is $nd \times D$ 307 *nd* specifying the covariance between *n* tracers at *d* locations 308 (depths). Assuming separability, Σ can be approximated by 309 a Kronecker product of two matrices

$$\Sigma = \Sigma_T \otimes C_S + \Sigma_M, \tag{4}$$

310 where Σ_T corresponds to the $n \times n$ cross-covariance matrix 311 of the tracers, and C_S is the $d \times d$ spatial correlation matrix 312 (in depth) respectively. Σ_M is the data measurement error 313 which we assume to be negligible compared to the other 314 errors because of the spatial aggregation of the data.

[20] The cross-covariance matrix Σ_T depends on n(n-1)316 2 cross-tracer correlation coefficients ρ_{ii} (since $\rho_{ii} = \rho_{ii}$), and 317 on residual standard deviations σ_i of the *n* individual tracers

$$\Sigma_{T} = \begin{bmatrix} \sigma_{1}^{2} & \sigma_{1}\sigma_{2}\rho_{12} & \dots & \sigma_{1}\sigma_{n}\rho_{1n} \\ \sigma_{2}\sigma_{1}\rho_{21} & \sigma_{2}^{2} & \dots & \sigma_{2}\sigma_{n}\rho_{2n} \\ \vdots & \vdots & \ddots & \vdots \\ \sigma_{n}\sigma_{1}\rho_{n1} & \dots & \dots & \sigma_{n}^{2} \end{bmatrix} . \tag{5}$$

318 [21] We model the spatial correlation C_S using a Gaussian 319 correlation function, a special case of the Matérn class of 320 covariance functions (see, for e.g., Stein [1999]). This 321 function decays with distance between locations d_i and d_i 322 with a correlation length scale λ , assumed to be the same for 323 all tracers

$$(C_S)_{ij} = \exp\left(-\frac{\left|d_i - d_j\right|^2}{\lambda^2}\right). \tag{6}$$

[22] Given the property of the Kronecker product (see, for 325 example, Lu and Zimmerman [2005]), the multivariate 326 normal likelihood function $L(v, \theta)$ becomes

$$L(O|\theta, \mu, \Sigma_T, C_S) = (2\pi)^{-N/2} \left(|\Sigma_T|^d |C_S|^n \right)^{-1/2} \cdot \exp\left[-\frac{1}{2} \tilde{r}^T \left(\Sigma_T^{-1} \otimes C_S^{-1} \right) \tilde{r} \right], \tag{7}$$

327 where N = nd is the total number of data points, and $\tilde{r} = [O_1 - O_1]$ 328 $M_1 - \mu_1, \ldots, O_n - M_n - \mu_n$ is the concatenated vector containing the misfit between the unbiased model predictions 329 and the corresponding observations for the considered tra- 330 cers. The Kronecker structure of equation (4) allows the $nd \times 331$ nd matrix Σ to be efficiently inverted by inverting the two 332 smaller matrices Σ_T ($n \times n$) and C_S ($d \times d$).

[23] Once the probability model has been specified in the 334 form of a likelihood function, the Bayes' theorem allows 335 inference about the posterior distribution of θ . The theorem 336 states that the posterior probability of the unknown parameters is proportional to their prior probability distribution, 338 multiplied by the likelihood of the data, according to

$$p(\theta, a, b, \sigma, \rho, \lambda | O) \propto L(O|\theta, a, b, \sigma, \rho, \lambda) p(\theta) p(a) \cdot p(b) p(\sigma) p(\rho) p(\lambda). \tag{8}$$

We draw 20,000 samples from the above posterior distri- 340 bution by a Markov chain Monte Carlo (MCMC) algorithm. 341 The MCMC algorithm jointly estimates the model para- 342 meters ($\theta = K_{bg}$, u_K_{SO}), 2n bias coefficients (a_i and b_i), n 343 standard deviations (σ_i) , n(n-1)/2 cross-tracer correlations 344 (ρ_{ii}) , and one correlation length (λ) . This is an improvement 345 upon the methodology of Schmittner et al. [2009] which 346 held all parameters but θ fixed at optimized values, and did 347 not consider the uncertainty in the other parameters. 348 Because the model output is only defined on a discrete grid 349 of values, the MCMC algorithm proposes discrete jumps for 350 the parameters θ during its random walk through parameter 351 space, and continuous moves for all other parameters.

[24] We choose a uniform prior $p(\theta)$ for the model parameters K_{bg} and $u_{-}K_{SO}$. For the correlation length we apply 354 the lognormal prior $\ln \lambda \sim N(5.5, 0.5^2)$, such that the logafithm of λ is normally distributed with mean 5.5 and standard deviation 0.5. This prior locates most of the probability 357 mass of the distribution between 0 and 600 meters. We use 358 normal priors for the bias parameters a_i and b_i , $p(a_i) = N(0,$ $(\sigma_i/\lambda)^2$) and $p(b_i) = N(0,\sigma_i^2)$. For the estimate of individual tracers distributions, where the cross-correlation matrix is a 361 scalar (i.e., $\Sigma = \sigma_1^2$), we use a Jeffreys prior $(p(\sigma_i) \propto 1/\Sigma)$. When the multitracer cross-covariance matrix is estimated, we specify an inverse Wishart prior distribution $\Sigma_T \sim IW(S,$ ν), with a diagonal scale matrix S = I and $\nu = 2n + 1$ degrees 365 of freedom. A diagonal scale matrix reduces spurious cor- 366 relations by penalizing tracer residuals which are not inde- 367 pendent of each other. Spurious correlation is not a problem 368 when the data dimension is large, but when the data are 369 sparse such a regularization procedure is prudent (see, for 370 instance, Barnard et al. [2000] or Chapter 19 of Gelman et 371 al. [2003], and references therein).

[25] Equation (8) provides the joint posterior probability 373 of both the model parameter and the bias and covariance 374 parameters. The marginal posterior probability of the model 375 parameter alone is obtained by integrating the joint posterior 376 over all other parameters,

$$p(\theta|O) = \int p(\theta, a, b, \sigma, \rho, \lambda|O) dadbd\sigma d\rho d\lambda. \tag{9}$$

[26] Since the posterior is estimated by MCMC sampling, 378 this posterior distribution of θ is easily obtained by simply 379 380 considering the θ samples while ignoring the samples for the 381 other parameters.

382 3. Results

383 3.1. Effect of Ocean Diapynal Diffusivity on the AMOC 384 Hindcasts and Spatial Fields

[27] In the adopted model the AMOC strength is positively 386 correlated with the parameters K_{bg} and $u_{\underline{K}SO}$ (Figure 1). K_{bg} 387 has a strong influence on the model hindcasts of the maxi-388 mum AMOC strength, while the AMOC sensitivity to 389 u_{KSO} is weaker. The range of AMOC strength varies from 390 about 5-23 Sv across all simulations. The inclusion of 391 enhanced upper Southern Ocean mixing ($u K_{SO} = 1$), can 392 increase the AMOC by a few Sverdrups, with more influence 393 at lower K_{bg} . Under the projected climate forcings, the 394 AMOC strength decreases in most cases, but it is more 395 sensitive (in absolute values) to the considered forcings for 396 higher diffusivity values. Due to the strong dependence of 397 the AMOC structure and behavior on the values of the 398 parameters K_{bg} and $u_{\underline{K}SO}$ in this model, a reduction in the 399 parametric uncertainty has the potential to improve AMOC 400 hindcast and projection in the model.

[28] The different parameter values result in different 402 hindcasts of ocean tracers such as CFC-11 (Figure 2) and $403 \Delta^{14}$ C (Figure 3), due to the different tracers advection and 404 diffusion rates in the model. Higher K_{bg} values result in 405 stronger vertical water exchange, increased deep water mass 406 formation, which carries water with higher tracer content 407 from the surface, and decreased vertical stratification in the 408 ocean. u_K_{SO} broadly produces the same effects of K_{bg} . 409 Nevertheless, u_K_{SO} impacts more heavily the lower K_{bg} 410 runs and the Southern Ocean stratification.

[29] Here we analyze the tracers concentrations as vertical 412 profiles of their averaged concentrations over the globe. We 413 consider three different observations, CFC-11, Δ^{14} C and T 414 (Figure 4, shown as an example for ENSEMBLE 1). In 415 general, the observations are contained by the model 416 ensemble spread, except for T in the deep ocean, which is 417 too cold in the model.

418 3.2. Uncertainty of the Statistical Inversion

[30] The inversion method uses the information contained 420 in the tracers to estimate the model parameter K_{bg} , taking 421 into account uncertainties in u K_{SO} . Key improvements 422 compared to Schmittner et al. [2009] are (1) the estimation 423 of the cross-correlation terms; (2) a more refined represen-424 tation of structured biases in the Likelihood function; and 425 (3) the consideration of the effects of the structural uncer-426 tainty (specifically about the implementation of mixing in 427 the SO). Here we demonstrate how these improvements 428 affect the joint posterior pdf of the model parameters. We 429 test the sensitivity of the method to the choice of the sta-430 tistical (or nuisance) parameters for the distribution of K_{bg} . 431 In this sensitivity test, we do not account for uncertainties in 432 the parameter u K_{SO} . Therefore, we only use outputs from 433 ENSEMBLE 1.

[31] For illustration, we use two tracers, Δ^{14} C and T, as 435 input for the statistical inversion. We calculate four inver-436 sion, which vary the number of statistical parameters to be 437 estimated. The structure of the errors differs from each other 438 by the representation of two main parameters, the bias and the cross correlation of the residuals between the model and 439 the observations. The bias term represents our guess of the 440 mean function of the residuals. We demonstrate the trade- 441 off between complexity of the bias-correction and the 442 covariance structure of the residuals in this simple sensi- 443 tivity study.

[32] Specifically, we analyze four different assumptions 445 about the structural error terms. First, we use a simple case 446 where the bias is constant and there is no residuals cross 447 correlation; second, we use a constant bias and estimate the 448 cross correlation; third we estimate a linear bias but no 449 residual cross correlation; and fourth, in which linear bias 450 and cross correlation are both estimated. To summarize the 451 experiments in the sensitivity study, we have (1) $\mu = b$, $\rho = 452$ 0, (2) $\mu = b$, $\rho = \hat{\rho}$, (3) $\mu = az + b$, $\rho = 0$, and (4) $\mu = az + b$, 453 $\rho = \hat{\rho}$. Note that the calibration also estimates standard 454 deviation, correlation length and the model parameter, as 455 described in section 2.3. Comparing all pdfs (Figure 5) we 456 see that for the individual pdfs the representation of the bias 457 term can be essential for the model parameter estimation. 458 When a more simplified bias ($\mu = b$) is applied (Figures 5a 459 and 5b), the pdfs in this example are displaced toward 460 higher K_{bg} values, and centered on 0.3 and 0.4 cm² s⁻¹. In 461 contrast, with the linear bias estimations, the mode of K_{bg} 462 pdf is centered around 0.15 and 0.2 cm² s⁻¹. For the cases 463 with linear bias (Cases c and d), the standard deviation of 464 the residuals of both tracers (Table 1) decrease in compar- 465 ison to the constant biases cases (Cases a and b). On the 466 other hand, the standard deviations of the residuals are not 467 influenced by the addition of cross-correlation parameters. 468

[33] The inclusion of the cross-correlation parameter im- 469 pacts the position of the joint posterior (black curves), and 470 its strength is closely related to the representation of the 471 bias. When the bias has a better representation, which is the 472 linear bias case here (Figures 5c and 5d), the cross-corre- 473 lation term has little influence on the joint pdf. A compar- 474 ison of the strength of the cross-correlation parameters 475 (Cases b and d in Table 1) shows that $\rho = 0.70$ when μ is 476 constant, and is much smaller $\rho = 0.40$ when μ is linear. 477 Comparing the posteriors of the Cases a and b (Figures 5a 478) and 5b), ρ can visibly change the posterior when the mean 479 function is less structured. Case b shows a counterintuitive 480 result where the posterior mode is distant from the modes of 481 the individual components (Figure 5b). This result indicates 482 that with a relatively poor representation of the mean (bias) 483 function, considering or neglecting the effects of this 484 residual cross-correlation can drastically change the K_{bg} 485 posterior estimate. This effect becomes less pronounced, as 486 the representation of the model bias term improves (e.g., 487 Figure 5b versus Figure 5c). As discussed by *Cressie* [1993] 488 (pp. 25), "What is one person's (spatial) covariance structure 489 may be another person's mean structure." In other words, 490 there is a trade-off between estimating a mean function for 491 the tracer residuals to account for structural model errors and 492 the magnitude of the residual cross correlation across the 493 considered sources of information.

3.3. Estimating the Uncertainty of Vertical Diffusivity 495

[34] The analysis so far illustrates how different tracers 496 observations can be combined to reduce uncertainty about 497 one mixing parameter (K_{bg}) . This reduction in parametric 498 uncertainty results, at least in the framework of the adopted 499

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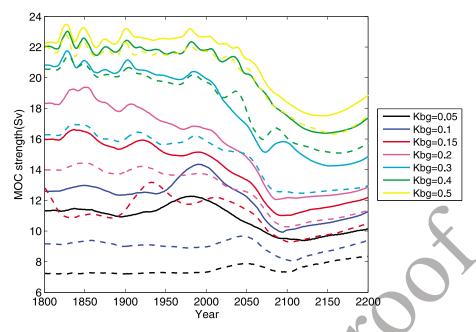


Figure 1. AMOC strength (Sv), defined as the maximum of the transport stream function, from years 1800 to 2200. Dashed lines are for the ENSEMBLE 1 ($u_K_{SO} = 0$); solid lines are for the ENSEMBLE 2 ($u_K_{SO} = 1$).

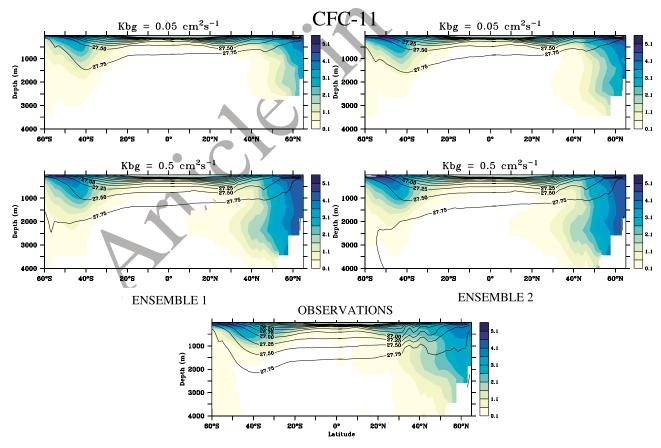


Figure 2. Zonal averages for the Atlantic Ocean of CFC-11 concentration in pmol/kg (color bars) and density anomalies in kg/m³ (contour lines) for the model with diffusivity of (top) $K_{bg} = 0.05$ and (middle) $K_{bg} = 0.5$. (left) ENSEMBLE 1 ($u_K_{SO} = 0$) and (right) ENSEMBLE 2 ($u_K_{SO} = 1$). (bottom) Observations from K_{ey} et al. [2004] and $L_{ocarnini}$ et al. [2006].

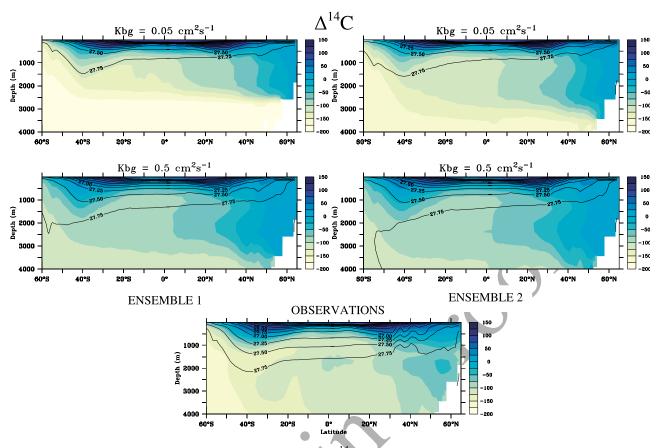


Figure 3. Zonal averages for the Atlantic Ocean of Δ^{14} C concentration in permil (color bars) and density anomalies in kg/m³ (contour lines) for the model with diffusivity of (top) $K_{bg} = 0.05$ and (middle) $K_{bg} = 0.5$. (left) ENSEMBLE 1 ($u_K_{SO} = 0$) and (right) ENSEMBLE 2 ($u_K_{SO} = 1$). (bottom) Observations from *Key et al.* [2004] and *Locarnini et al.* [2006].

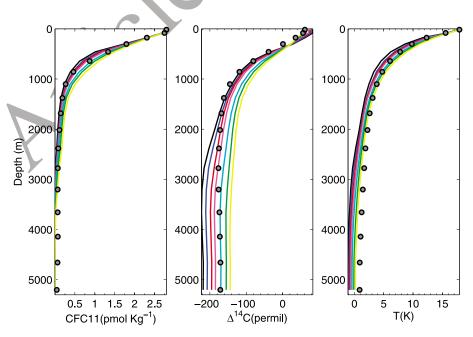


Figure 4. Global averaged profiles of CFC-11 [Key et al., 2004], Δ^{14} C [Key et al., 2004], and T [Locarnini et al., 2006] for the observations (gray dots) and model ENSEMBLE 1 (colored lines). The legend for the model K_{bg} values is the same as in Figure 1.

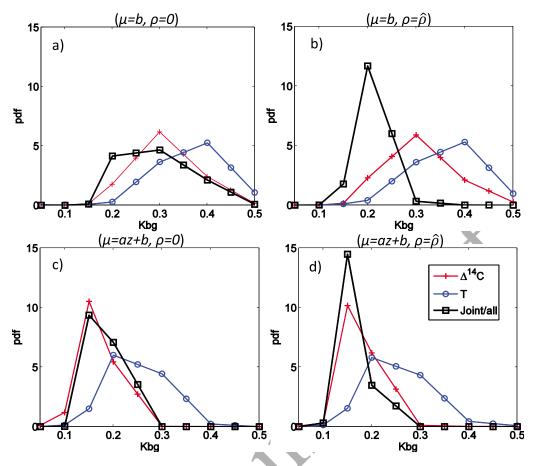


Figure 5. Sensitivity of the model parameter estimation to different treatments of structural model errors. Shown are the posterior probability density function of Δ^{14} C (red lines with crosses) and T (blue lines with circles), and the joint posterior using both observations (black line with squares). The frames are for the cases discussed in the text: (a) $[\mu = b, \rho = 0]$, (b) $[\mu = b, \rho = \hat{\rho}]$, (c) $[\mu = az + b, \rho = 0]$, and (d) $[\mu = az + b, \rho = \hat{\rho}]$.

500 model, in a reduction of the prediction uncertainty about the 501 AMOC. Of course, there are several caveats associated with 502 structural errors and other neglected uncertainties in this 503 study. We return to this issue in section 4. In this section we 504 illustrate how this information can potentially be used to 505 reduce uncertainties in two model parameters and improve 506 model hindcasts and projections of the AMOC. Here the 507 inversion uses our best estimate of the model bias term 508 (linear), and accounts for cross-tracer correlation. We make 509 three inversions (Figure 6), one to estimate K_{bg} for the 510 ENSEMBLE 1, a second to estimate K_{bg} for the ENSEM-511 BLE 2, and a third inversion which uses information from

t1.8 t1.9 both ensembles to generate probability distributions for K_{bg} 512 and u_K_{SO} in a Bayesian model average fashion. 513

[35] Information from the three considered tracers, CFC- 514 11, T and Δ^{14} C, is introduced in the statistical inversion for 515 the estimation of uncertainties in the model parameters. For 516 comparison, we also show in Figure 6 the K_{bg} pdf obtained 517 using the climatological AMOC observations. The K_{bg} pdf 518 is derived from estimate of the climatological AMOC 519 strength of *Lumpkin and Speer* [2003] by assimilating a 520 single data point assuming a normally distributed error. In 521 principle, the model could be calibrated with both the ocean 522 tracers and AMOC strength data by using the derived 523

t1.1 **Table 1.** Properties of the Statistical Distributions of the Sensitivity Test for the Best K_{bg}^{a}

		$\frac{\text{Mode}}{(\text{cm}^2 \text{ s}^{-1})}$		Bias (a,b)		σ		Cross Correlation	Mode of
t1.3	Experiment	Δ^{14} C	T	$\Delta^{14}\mathrm{C}$	T	Δ^{14} C	T	at Best K_{bg}	Posterior
t1.4	Case a	0.3	0.4	(-14.0,0)	(0.45,0)	12.5	0.6	_	0.3
t1.5	Case b	0.3	0.4	(-14.0,0)	(0.45,0)	12.5	0.6	0.70	0.2
t1.6	Case c	0.15	0.2	(-16.1,9e-3)	(0.22, 3.3e-4)	7.7	0.28	_	0.15
t1.7	Case d	0.15	0.2	(-16.1,9e-3)	(0.22,3.2e-4)	7.7	0.28	0.40	0.15

^aMode, bias ($\mu = az + b$), standard deviation, and cross-correlation of residuals for Δ^{14} C and T and mode of the posterior (joint distribution considering all tracers information).

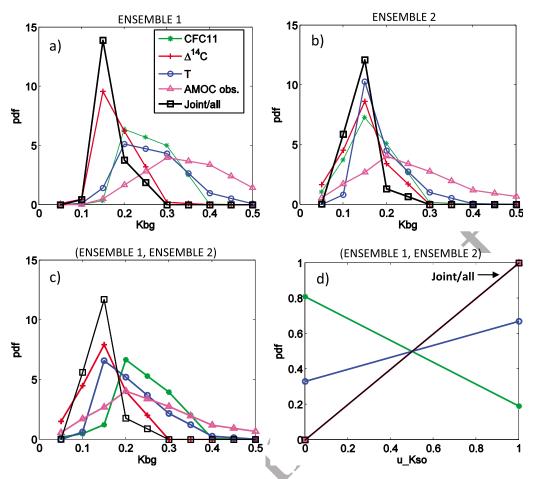


Figure 6. Posterior probability density function of the model parameters for all considered sources of information, the joint posterior using all available information from observations (black line with squares). The climatological AMOC estimate of Lumpkin and Speer [2003] is plotted for comparison (pink line with triangles). The K_{bg} estimates are for (a) ENSEMBLE 1, (b) ENSEMBLE 2, (c) ENSEMBLE 1 and ENSEMBLE 2, and (d) the $u_{-}K_{SO}$ estimate is for ENSEMBLE 1 and ENSEMBLE 2.

524 AMOC pdf as a prior for K_{bg} . However, this would neglect 525 potential correlations between ocean tracer and AMOC 526 strength residual errors. As a proper treatment of AMOC/ 527 tracer correlations is beyond the scope of this work, we 528 present the AMOC-derived pdf just for comparison, without 529 assimilating it in the joint posterior pdf.

[36] The tracers distributions of both ensembles show sim-531 ilar behavior. Nevertheless, the ENSEMBLE 1 (Figure 6a) has 532 in general higher K_{bg} modes in comparison to ENSEMBLE 2 533 (Figure 6b). This result shows that the additional mixing over 534 the upper Southern Ocean increases the overall magnitude 535 of K_v , without changing K_{bg} , and tends to intensify the AMOC. 536 The posterior pdf for K_{bg} , obtained by assimilating observa-537 tions of the AMOC strength only (lines with triangles), is 538 also displaced to lower values in ENSEMBLE 2, because 539 ENSEMBLE 2 has stronger AMOC values for the same K_{bg} 540 (Figure 1).

[37] When information from both ensembles are added 542 together (Figures 6c and 6d), the ENSEMBLE 2 dominates 543 the Markov chain for Δ^{14} C and T, with probabilities of 544 100% and 65% for ENSEMBLE 2, respectively. Con-545 versely, CFC-11 has 80% probability of happening 546 ENSEMBLE 1 (Figure 6d). The joint posterior of all tracers

encompassing the two ensembles (Figure 6c) is entirely 547 described by ENSEMBLE 2; therefore, the posteriors in 548 Figures 6b and 6c are practically identical.

[38] When all the two model parameters are assimilated 550 jointly (Figure 6c), the considered sources of information 551 have rather different skill in improving K_{bg} estimates and 552 AMOC predictions (see Table 2 for the properties of the 553 statistical distributions). Δ^{14} C has the highest information 554 content with respect to improving K_{bg} estimates, its poste- 555 rior 95% credible interval (CI) is the tightest (0.21 cm² s⁻¹) 556 in comparison to the other tracers. CFC-11 comes in second, 557 with a 95% CI of 0.24 cm² s⁻¹, and T comes last with the 558 largest CI of 0.26 cm² s⁻¹.

[39] Combining the information of the three considered 560 tracers (line with squares in Figure 6c), favors K_{bg} values in 561 the lower part of the considered range, from $0.1 \text{ to } 0.2 \text{ cm}^2 \text{ s}^{-1}$. 562 Note that the joined probability density function is narrower 563 than each individual pdf indicating an advantage of using 564 multiple tracer observations in reducing the parameter 565 uncertainty.

[40] As discussed in previous studies [e.g., Schmittner et al. 567 2009], the K_{bg} value in a coarse resolution ocean model 568 represents the effects of background diffusivity combined 569 t2.12

t2.13

t2.14

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Table 2. Properties of the Statistical Distributions of K_{bg} (Figure 6c) t2.1 for Each Considered Sources of Information, the Posterior (Joint t2.2 t2.3 Distribution Considering All Tracers Information), and the Climatological AMOC Estimate^a t2.4

					Cross Correlation at Best K_{bg}		
t2.6	Observation	Mode	Mean	95% CI	Δ^{14} C	CFC-11	T
t2.7	Δ^{14} C	0.15	0.15	0.22	1	0.06	0.38
t2.8	CFC-11	0.20	0.23	0.26	_	1	0.02
t2.9	T	0.15	0.18	0.26	_	_	1
t2.10	Posterior	0.15	0.16	0.17	_	_	_
t2.11	Climatological sAMOC	0.20	0.20	0.42	-	_	_

^aMode, mean, and 95% credible interval (CI, in cm² s⁻¹). Climatological AMOC estimate from Lumpkin and Speer [2003]. Also shown are the cross-tracer correlation at the best K_{bg} value estimated in the joint posterior.

570 with subgridscale diffusivity (i.e., a model shortcoming). 571 Another shortcoming for coarse z coordinate ocean models is 572 the numerical diffusivity (Veronis effect), which can generate 573 spurious diapycnal diffusion, especially in long climate si-574 mulations, in western boundary regions and regions where 575 the isoneutral slope is large [Griffies et al., 1998, 2000]. 576 Hence, even if our model-based estimate does not represent 577 directly the observational estimate of pelagic diffusivity of $578 \ 0.1 \ \text{cm}^2 \ \text{s}^{-1}$ [Ledwell et al., 1993], they appear to be more 579 consistent when we improve on the parameterization of 580 regional mixing in the model.

581 3.4. AMOC Projections

[41] The joint posterior K_{bg} and $u_{\underline{K}SO}$ estimates 583 (Figures 6c and 6d) can be used to derive model projections of 584 the AMOC in 2100 and 2200 (Figure 7). The model hindcast 585 for the maximum AMOC strength in 2000 is about 15-586 15.5 Sv. In 2100, the expected strength for the AMOC in 587 this model is about 11 Sv. In 2200 the AMOC shows a 588 slight strengthening relative to the 2100 conditions with an 589 expected value of roughly 12 Sv.

[42] The K_{bg} and u_K_{SO} estimates suggest an AMOC 591 hindcast for the year 2000 (Figure 7) that is about 2 Sv 592 weaker than the climatological AMOC estimates of Lumpkin 593 and Speer [2003]. The inclusion of the parameter $u_{-}K_{SO}$ in 594 the analysis reduces significantly the discrepancy of the 595 AMOC estimates relative to the K_{bg} (Figure 6c). Other sys-596 tematic model bias(es), such as too weak buoyancy forcing 597 (e.g., from errors in the simulation of the atmospheric hydro-598 logical cycle and surface freshwater fluxes) can compromise 599 the estimates of the current and projected AMOC strength 600 for the Uvic model. Further discussion and implications are 601 described in section 4.

602 4. Caveats

[43] Our results are subject to many caveats. These caveats 604 point to potentially fruitful research directions. In the statis-605 tical part, we consider only highly aggregated data. Basin-606 wide zonal averages could, for example, provide potentially 607 useful information on where the model performs better. In 608 the projection part, other model parameters, such as those 609 affecting the response of the ocean-atmosphere coupled 610 system, for example, the hydrological cycle [Saenko and

Weaver, 2004], climate sensitivity or sensitivity of climate 611 to aerosol concentrations, [cf. Tomassini et al., 2007; Forest 612 et al., 2002], are also highly uncertain, and can impact 613 (probably widen) probabilistic AMOC projections and 614 should be considered. In addition, the atmospheric model in 615 UVic is rather simplified, and neglects important ocean- 616 atmosphere feedbacks.

[44] UVic does not use flux correction. Freshwater flux 618 correction is known to improve the salinity and stratification 619 in ocean models [Sorensen et al., 2001], and can be used to 620 improve projections and hindcasts.

[45] In the hindcasts part, other parameters linked to both 622 diapycnal and isopycnal mixing may affect the structure of 623 the AMOC. Nevertheless, according to Jayne [2009], tidal 624 mixing parameters in the [St. Laurent et al., 2002] param- 625 eterization have relatively low impact on the strength of the 626 AMOC, and that upper-ocean wind-driven mixing may 627 have a much stronger impact.

[46] We show how including regional aspects of vertical 629 mixing can improve the representation of the AMOC. The 630 model parameters uncertainties need to be estimated 631 together as performed here, since addition of new para- 632 meters can change the structure of the other calibrations. 633 Jayne [2009] describes, "this is the typical conundrum: it is difficult to assess whether any of the given parameterizations improve the model since comparing to observational 636 metrics may obscure compensating errors in different para-637 meterizations."

Conclusion 639

[47] We develop and apply a computationally efficient 640 and statistically sound method to rank and quantify the skill 641 of different sources of information to reduce the uncertainty 642 about ocean model parameters and the resulting climate 643 predictions. We improve on previous work by (1) refining 644 the estimation of errors in the model structure, (2) including 645 several ocean tracers and two model parameters at once in a 646 computationally efficient fashion, and (3) quantifying and 647 ranking the skill of different sources of information to 648

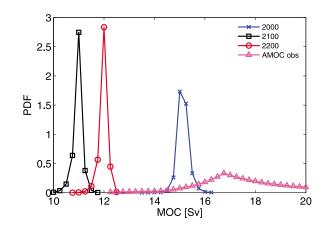


Figure 7. Joint posterior probability density function of model projections of the maximum AMOC strength in the years 2000, 2100, and 2200 using information from the Δ^{14} C, CFC-11, and T observations. The climatological AMOC estimate of Lumpkin and Speer [2003] is added for comparison (pink line with triangles).

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- 649 reduce the uncertainty about a model parameter. Subject to 650 the aforementioned caveats, we show how Δ^{14} C, CFC-11, 651 and T together sharpen the estimates of K_{bg} by 40% and 652 improve AMOC projections in the UVic model.
- [48] The K_{bg} derived from individual observations (i.e., $654 \Delta^{14}$ C, CFC- $1\overline{1}$, T) are broadly consistent, but show slight 655 discrepancies that we attribute predominantly to structural 656 model errors. Of the considered observations, Δ^{14} C has the 657 highest skill in reducing uncertainties in AMOC projections, 658 but it is also the most distant from the pdf observational 659 derived AMOC estimates. Δ^{14} C is followed (in decreasing 660 skill of being able to reduce K_{bg} uncertainty) by CFC-11 661 and T. The second parameter analyzed in this work, u_{KSO} 662 improved the representations of C14 and T in the model, and 663 improves the representation of the AMOC strength.
- [49] AMOC projections show a reduction of the maxi-665 mum of the joint posterior in 2100 by roughly 25% (3.5 Sv). 666 Perhaps both surprisingly and encouraging, the pdfs of K_{bg} 667 estimated in this study are quite similar among the consid-668 ered ocean tracers and the two ensembles analyzed, which 669 have different representations of the upper Southern 670 Oceanmixing and AMOC. This convergence of K_{bg} esti-671 mates based on different sources of information and para-672 meterizations suggest that K_{bg} can be robustly estimated 673 from the oceanic tracers studied here.
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681 References

- 682 Baehr, J., D. McInerney, K. Keller, and J. Marotzke (2008), Optimization of an observing system design for the North Atlantic Meridional Over-
- turning Circulation, *J. Atmos. Oceanic Technol.*, 25, 625–634. 885 Barnard, J., R. McCulloch, and X. Meng (2000), Modeling covariance 686 matrices in terms of standard deviations and correlations with application to shrinkage, Stat. Sinica, 10, 1281-1311.
- Bhat, K. S., M. Haran, R. Tonkonojenkov, and K. Keller (2009), Inferring 688 689 likelihoods and climate system characteristics from climate models and 690 multiple tracers, technical report, Dep. of Stat., Penn. State Univ., Univer-691 sity Park, Pa.
- 692 Broecker, W. S., and T. H. Peng (1987), Gas exchange rates between air 693 and sea, *Tellus*, *26*, 21–35.
- 694 Bryan, K., and L. J. Lewis (1979), A water mass model of the world ocean, 695 *J. Geophys. Res.*, 84(C5), 2503–2517, doi:10.1029/JC084iC05p02503.
- 696 Cressie, N. A. (1993), Statistics for Spatial Data, 900 pp., Wiley, New 697
- 698 Doney, S. C., et al. (2004), Evaluating global ocean carbon models: The 699 importance of realistic physics, Global Biogeochem. Cycles, 18, 700GB3017, doi:10.1029/2003GB002150.
- 701 England, M. H. (1993), Representing the global scale water masses in ocean general circulation models, J. Phys. Oceanogr., 23, 1523-1552.
- 703 Ferrari, R., J. C. McWilliams, V. M. Canuto, and M. Dubovikov (2008), 704 Parameterization of eddy fluxes near oceanic boundaries, J. Clim., 21, 7052770-2789, doi:10.1175/2007JCLI1510.1.
- 706 Forest, C. E., P. H. Stone, A. P. Sokolov, M. R. Allen, and M. D. Webster (2002), Quantifying uncertainties in climate system properties with the use of recent climate observations, Science, 295, 113-117, doi:10.1126/ 709
- science.1064419. 710 Gao, Y., H. Drange, and M. Bentsen (2003), Effects of diapycnal and isopycnal mixing on the ventilation of CFCs in the North Atlantic in an isopycnic
- coordinate OGCM, Tellus Ser. B, 55, 837-854. 713 Gelman, A., J. B. Carlin, H. S. Stern, and D. B. Rubin (2003), Bayesian
- Data Analysis, 2nd ed., CRC Press, London.
- 715 Gent, P. R., and J. C. McWilliams (1990), Isopycnal mixing in ocean circulation models, J. Phys. Oceanogr., 20, 150-155.

- Gnanadesikan, A. (1999), A simple predictive model for the structure of the 717 oceanic pycnocline, Science, 283, 2077-2079.
- Griffies, S. M., A. Gnanadesikan, R. C. Pacanowski, V. D. Larichev, J. K. Dukowicz, and R. D. Smith (1998), Isoneutral diffusion in a z coordinate ocean model, J. Phys. Oceanogr., 28, 805-830.
- Griffies, S. M., R. C. Pacanowski, and R. W. Hallberg (2000), Spurious diapycnal mixing associated with advection in a z-coordinate ocean model, Mon. Weather Rev., 128(3), 538-564.
- Hansen, J., and M. Sato (2004), Greenhouse gas growth rates, Proc. Natl. Acad. Sci. U.S.A., 101, 16,109-16,114.
- Ito, T., J. Marshall, and M. Follows (2004), What controls the uptake of transient tracers in the Southern Ocean?, Global Biogeochem. Cycles, 18, GB2021, doi:10.1029/2003GB002103.
- 730Jayne, S. R. (2009), The impact of abyssal mixing parameterizations in an ocean general circulation model, J. Phys. Oceanogr., 39, 1756-1775. 731732
- Keller, K., K. Tan, F. M. M. Morel, and D. F. Bradford (2000), Preserving the ocean circulation: Implications for climate policy, Clim. Change, 47,
- Keller, K., C. Deutsch, M. G. Hall, and D. F. Bradford (2007), Early detection of changes in the North Atlantic Meridional Overturning Circulation: Implications for the design of ocean observation systems, J. Clim., 20,
- Key, R. M., A. Kozyr, C. L. Sabine, K. Lee, R. Wanninkhof, J. Bullister, R. A. Feely, F. Millero, C. Mordy, and T.-H. Peng (2004), A global ocean carbon climatology: Results from Global Data Analysis Project (GLODAP), Global Biogeochem. Cycles, 18, GB4031, doi:10.1029/ 2004GB002247.
- Koch, D., D. Jacob, I. Tegen, D. Rind, and M. Chin (1999), Tropospheric sulfur simulation and sulfate direct radiative forcing in the Goddard Institute for Space Studies general circulation model, J. Geophys. Res., 104(D19), 23,799-23,822, doi:10.1029/1999JD900248.
- Large, W. G., J. C. McWilliams, and S. C. Doney (2004), Oceanic vertical mixing: A review and a model with nonlocal boundary layer parameterisation, Rev. Geophys., 32, 363-403.
- Ledwell, J. R., A. J. Watson, and C. S. Law (1993), Evidence for slow mixing across the pycnocline from an open-ocean tracer-release experiment, Nature, 364, 701-703.
- ocarnini, R. A., A. V. Mishonov, J. Antonov, T. Boyer, and H. E. Garcia (2006), World Ocean Atlas 2005, vol. 1, Temperature, NOAA Atlas 754 755NESDIS, vol. 61, 182 pp., NOAA, Silver Spring, Md. 756757
- Lu, N., and D. L. Zimmerman (2005), The likelihood ratio test for a separable covariance matrix, Stat. Prob. Lett., 73(4), 449-457. 759
- Lumpkin, R., and K. Speer (2003), Large-scale vertical and horizontal circulation in the North Atlantic Ocean, J. Phys. Oceanogr., 33, 1902–1920.
- Marotzke, J. (1997), Boundary mixing and the dynamics of threedimensional thermohaline circulation, J. Phys. Oceanogr., 27, 1713–1728.
- McManus, J. F., R. Francois, J. M. Gherardi, L. D. Keigwin, and S. Brown-Leger (2004), Collapse and rapid resumption of Atlantic meridional circulation linked to deglacial climate changes, Nature, 428, 834-837.
- Meehl, G. H., et al. (2007), Global climate projections, in Climate Change 2007: The Physical Science Basis: Contribution of Working Group I to the Fourth Assessment Report of the Intergovernmental Panel on Climate Change, edited by S. Solomon et al., pp. 747-846, Cambridge Univ. Press, Cambridge, U. K.
- Metropolis, N., A. W. Rosenbluth, M. N. Rosenbluth, A. H. Teller, and E. Teller (1953), Equation of state calculations by fast computing machines, J. Chem. Phys., 21, 1087-1092.
- Moum, J. N., and W. D. Smyth (2001), Upper ocean mixing processes, in Encyclopedia of Ocean Sciences, edited by J. H. Steele, K. K. Turekian, and S. A. Thorpe, pp. 3093-3100, Academic, New York.
- Munk, W., and C. Wunsch (1998), Abyssal recipes: Part II. Energetics of
- tidal and wind mixing, *Deep Sea Res. Part I, 13*, 1977–2010. Naveira Garabato, A. C., K. L. Polzin, B. A. King, K. J. Heywood, and M. Visbeck (2004), Widespread intense turbulent mixing in the Southern Ocean, Science, 303, 210-213.
- Pacanowski, R., and S. G. H. Philander (1981), Parameterization of vertical mixing in numerical models of tropical oceans, J. Phys. Oceanogr., 11, 1443-1451
- Pacanowski, R. C. (1995), MOM 2 documentation: Users guide and reference manual, ver. 1.0, Ocean Group Tech. Rep. 3, Geophys. Fluid Dyn. Lab., Princeton, N. J.
- Patwardhan, A., et al. (2007), Assessing key vulnerabilities and the risk from climate change, in Climate Change 2007: Impacts, Adaptation and Vulnerability: Contribution of Working Group II to the Fourth Assessment Report of the Intergovernmental Panel on Climate Change, edited by M. L. Parry et al., chap. 19, pp. 779-810, Cambridge Univ. Press, Cambridge, U. K.
- Saenko, O. A., and A. J. Weaver (2004), What drives heat transport in the Atlantic: Sensitivity to mechanical energy supply and buoyancy forcing 795

841

842

843

- 796 in the Southern Ocean, *Geophys. Res. Lett.*, *31*, L20305, doi:10.1029/797 2004GL020671.
- 798 Sato, M., J. E. Hansen, M. P. McCormick, and J. B. Pollack (1993), Strato rope spheric aerosol optical depth, 1850–1990, J. Geophys. Res., 98(D12),
 22,987–22,994, doi:10.1029/93JD02553.
- 801 Schmittner, A., and A. J. Weaver (2001), Dependence of multiple cli-802 mate states on ocean mixing parameters, *Geophys. Res. Lett.*, 28(6), 803 1027–1030, doi:10.1029/2000GL012410.
- 804 Schmittner, A., A. Oschlies, X. Giraud, M. Eby, and H. L. Simmons
 805 (2005), A global model of the marine ecosystem for long-term simula806 tions: Sensitivity to ocean mixing, buoyancy forcing, particle sinking,
 807 and dissolved organic matter cycling, Global Biogeochem. Cycles, 19,
 808 GB3004, doi:10.1029/2004GB002283.
- 809 Schmittner, A., N. M. Urban, K. Keller, and D. Matthews (2009), Using
 810 tracer observations to reduce the uncertainty of ocean diapycnal mixing
 811 and climate carbon-cycle projections, *Global Biogeochem. Cycles*, 23,
 812 GB4009, doi:10.1029/2008GB003421.
- 813 Simmons, H. L., S. Jayne, L. St. Laurent, and A. Weaver (2004), Tidally 814 driven mixing in a numerical model of the ocean general circulation, 815 *Ocean Modell.*, 6, 245–263, doi:10.1016/S1463-5003(03)00011-8.
- 816 Smyth, W. D., and J. N. Moum (2001), Three-dimensional turbulence, in
 817 Encyclopedia of Ocean Sciences, edited by J. H. Steele, K. K. Turekian,
 818 and S. A. Thorpe, pp. 18–25, Academic, London.
- 819 Sokolov, A., C. Wang, G. Holian, P. Stone, and R. Prinn (1998), Uncer-820 tainty in the oceanic heat and carbon uptake and its impact on climate 821 projections, *Geophys. Res. Lett.*, 25(19), 3603–3606, doi:10.1029/ 822 98GL02696.
- 823 Sorensen, J., J. Ribbe, and G. Shaffer (2001), Antarctic intermediate water
 824 mass formation in ocean general circulation models, *J. Phys. Oceanogr.*,
 825 31, 3295–3311.
- 826 Sriver, R. L., M. Goes, M. E. Mann, and K. Keller (2010), Climate
 827 response to tropical cyclone-induced ocean mixing in an Earth system
 828 model of intermediate complexity, *J. Geophys. Res.*, doi:10.1029/
 829 2010JC006106, in press.
- 830 Stein, M. L. (1999), *Interpolation of Spatial Data: Some Theory for Kriging*, 831 247 pp., Springer, New York.

- St. Laurent, L., H. L. Simmons, and S. R. Jayne (2002), Estimating tidally driven mixing in the deep ocean, *Geophys. Res. Lett.*, 29(23), 2106, doi:10.1029/2002GL015633.
- Stuiver, M., and H. A Polach (1977), Discussion: Reporting of ¹⁴C data, 835 *Radiocarbon*, 19(3), 355–363.
- Toggweiler, J. R., K. Dixon, and K. Bryan (1989), Simulations of radiocarbon in a coarse-resolution, world ocean model: 2. Distributions of 838 bomb-produced ¹⁴C, *J. Geophys. Res.*, 94(C6), 8243–8264, 839 doi:10.1029/JC094iC06p08243.
- Tomassini, L., P. Reichert, R. Knutti, T. F. Stocker, and M. Borsuk (2007), Robust Bayesian uncertainty analysis of climate system properties using Markov chain Monte Carlo methods, J. Clim., 20, 1239–1254.
- Weaver, A. J., et al. (2001), The UVic Earth system climate model: Model 844 description, climatology, and applications to past, present and future climates, *Atmos. Ocean*, 39(4), 361–428.
- Wunsch, C., and R. Ferrari (2004), Vertical mixing, energy, and the general 847 circulation of the oceans, *Ann. Rev. Fluid Mech.*, 36, 281–314.
- Zickfeld, K., A. Levermann, M. Morgan, T. Kuhlbrodt, S. Rahmstorf, and D. Keith (2007), Expert judgements on the response of the Atlantic Meridional Overturning Circulation to climate change, *Clim. Change*, 851 82, 235–265.
- Zickfeld, K., M. Eby, and A. J. Weaver (2008), Carbon cycle feedbacks of changes in the Atlantic Meridional Overturning Circulation under future atmospheric CO₂, Global Biogeochem. Cycles, 22, GB3024, 855 doi:10.1029/2007GB003118.
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