

**INVESTIGATION OF A HIPPED PLATE ROOF STRUCTURE**

**by**

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# INVESTIGATION OF A HIPPED PLATE ROOF STRUCTURE

## PART I

### INTRODUCTION

Folded plate thin shells for roofs of large buildings are defined as structures in which the loads are sustained by thin folded flat plates. These folded plate structures, or as they are sometimes called, prismatic shells or hipped plates, have an economic and aesthetic advantage over certain other methods of roofing large buildings. This type has an advantage over thin curved shells in that it is easier to form, fabricate and design.

This paper deals with a "three-segment (horizontal projection)" plywood hipped plate roof simply supported by two end diaphragms. Although this type of structure takes advantage of the effective use of the monolithic properties of the material such as reinforced concrete, the writer attempted to extend both theoretical and experimental investigations to the plywood structure which is not homogeneous, nor monolithic.

The first folded plate structure was constructed in Germany in 1925 for use as coal bunkers where the ratio of the span length to the width of plate is large.

G. Ehlers and H. Craemer published the first design theory in 1930, in which the rigidity of the joints was entirely neglected. In 1932 E. Gruber published a paper in which he considered the effect of the rigidity of the joints, the connecting moments along the common edges of the plates, and the effect of the relative displacements between the joints (4, p. 744). The practice of this type of roof started in America as early as 1935 when L. H. Nishkan designed and built a warehouse in San Francisco using hipped plate theory. Since then, however, the folded plate construction had been ignored preceding World War II. After the War many investigations have been conducted. Interest in this type of roof has increased considerably since the first paper on the design theory was published in this country by Winter and Pei in 1947. They adopted the method of analysis in the European countries to the American practice and developed a distribution method (11, p. 506). Among many investigations, the most valuable contribution to this field was done by I. Gaafar in 1954. Gaafar did both theoretical and experimental work on an aluminum model of a three-segment hipped plate roof. In his analysis, Gaafar included joint

displacements as well as previous fundamental hipped plate theories. It was not new to consider the joint displacements, but Gaafar simplified it tremendously. However, he neglected the fact that the longitudinal edges of plates are on the elastic supports, such that a part of the uniformly distributed load in longitudinal direction is translated transversely by the slab action due to the displacement along the longitudinal edges of the plate. In a relaxation method for the corrected longitudinal stresses, it is more accurate and simpler to solve simultaneous equations than to use a moment distribution method.

In designing the folded plate, the longitudinal bending of a plate (out of its plane) takes an important role in addition to the transverse slab action and longitudinal plate action. This was witnessed by the difference in the longitudinal strain gage readings between inner and outer surface of the model, especially in the edge beam. This type of bending would be a critical factor in designing reinforced concrete hipped plate structures.

## PART II

## THEORY

A. Ordinary Beam Approach

Assuming no relative displacement of joint and no relative rotations between the adjacent plates, the structure is treated as a box beam. In this type of approach the loads are sustained by the beams at all joints along the edges of the plate.

B. Assumptions Made in Analyzing the Simply Supported Hipped Plate Structure

1. The material is uniform and elastic.
2. The cross section of the plate is rectangular shape.
3. All joints are rigid and torsional resistances or effects on the plate are neglected.
4. Stresses vary linearly across the width of each plate.
5. The ratio of length to the width of the plate is large, such that the plate can be treated as a one-way slab.

### C. Hipped Plate Theory

Since a part of the surface load is transmitted to the longitudinal edges of the plate, because of the difficulty in expressing such a loading condition, the corrective surface load can be approximated by the sum of partial loads which produces the sinusoidal elastic line for a simply supported roof structure. This complies with the previous assumptions in which the elastic line of a beam carrying uniform or concentrated load is a half sine wave. The intensity of the sum of partial load is defined as  $A_n \sin \frac{n\pi x}{l}$ , in which  $n$  is any integer. This trigonometric function is defined as a series of the form

$$\begin{aligned} & \frac{a_0}{2} + a_1 \cos \frac{\pi x}{l} + a_2 \cos \frac{2\pi x}{l} + \dots + a_n \cos \frac{n\pi x}{l} \\ & + b_1 \sin \frac{\pi x}{l} + b_2 \sin \frac{2\pi x}{l} + \dots + b_n \sin \frac{n\pi x}{l} . \end{aligned}$$

Expand this trigonometric function defined by the relations,



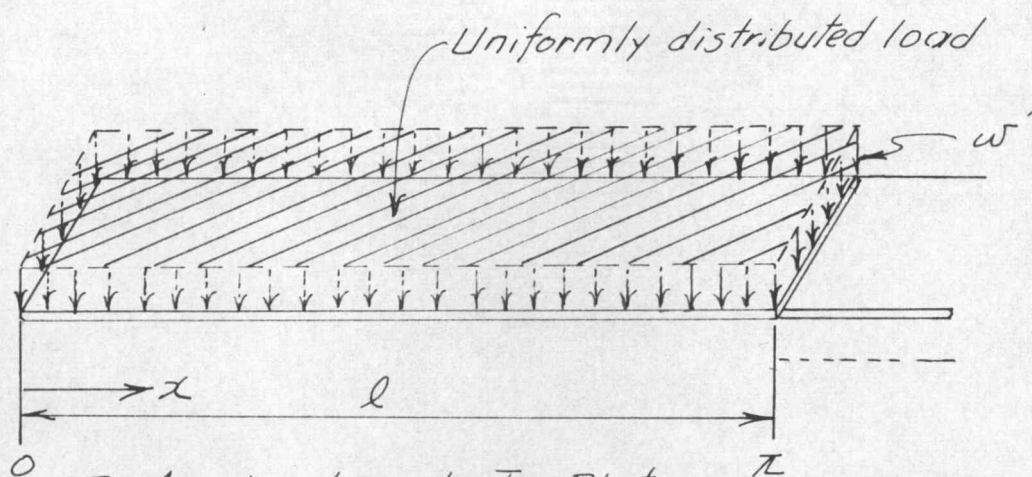


Fig.1 — Load on the Top Plate

$$f(x) = w', \quad 0 < x < \pi \text{ or } l,$$

$$f(x) = -w', \quad l \text{ or } \pi < x < 2\pi \text{ or } 2l,$$

in Fourier series.

From the Fourier formulas (5, p. 138),

$$a_0 = \frac{1}{\pi} \int_c^{c+2\pi} f(x) dx = \frac{1}{\pi} \int_0^{2\pi} f(x) dx = \frac{1}{\pi} \int_0^{2l} f(x) dx = 0$$

$$a_n = \frac{1}{\pi} \int_0^{2\pi} f(x) \cos \frac{n\pi x}{l} dx = 0$$

$$b_n = \frac{1}{\pi} \int_0^{2\pi} f(x) \sin \frac{n\pi x}{l} dx = \frac{1}{\pi} \int_0^l w' \sin \frac{n\pi x}{l} dx$$

$$+ \frac{1}{\pi} \int_l^{2l} (-w') \sin \frac{n\pi x}{l} dx = \frac{2w'}{n\pi} [1 - \cos n\pi]$$

$$\begin{aligned}
 f(x) = p &= \sum_{n=1}^{\infty} b_n \sin \frac{n\pi x}{l} \\
 &= \frac{4w'}{\pi} \left[ \sin \frac{\pi x}{l} + \frac{1}{3} \sin \frac{3\pi x}{l} + \dots \right] \\
 &= \frac{4w'}{\pi} \sum_{n=1,3,5,\dots} \frac{1}{n} \sin \frac{n\pi x}{l} \quad \text{-----} (1)
 \end{aligned}$$

For simplicity and sufficient accuracy for preliminary design purpose, the first term of Eq. (1) is used. Line load along the longitudinal direction, then, is

$$p = \frac{4w'}{\pi} \sin \frac{\pi x}{l} \quad \text{-----} (1')$$

Load in the transverse direction varies uniformly. The sections, transversely cut as a continuous one-way slab supported by imaginary reactions at the joints of the longitudinal plates, are analyzed. Analyzing a continuous one-way slab, any type of loads can be transformed into the loads acting at the joints. These joint loads are resolved into two components parallel to the two adjacent plates. The components of the joint loads cause the independent bending in each plate which acts as a deep beam simply supported by end diaphragms. Unequal longitudinal edge stresses at the joints do not comply with the continuity requirements. Consequently, the shear stresses are bound to act along the common edges in order to prevent



the relative movements of the adjacent plates. Creating the imaginary two equal and opposite forces at the common edges, the corrected longitudinal stresses are found.

The relative displacement of joints are derived from the plate deflections which cause the plate to rotate. Again analyzing a continuous one-way slab, plate loads are found. Then the process for unyielding supports is repeated. The longitudinal plate stresses (out of plane) are found by using the deflection perpendicular to the plate.

#### D. Analytical Criteria for Design

Stresses to be found for designing the hipped plate roof structure are divided into three stages:

1. the longitudinal stresses.
2. the transverse stresses.
3. the longitudinal (out of plane) stresses.

### PART III

#### DESIGN OF A PLYWOOD MODEL

##### A. Material Descriptions

Douglas fir plywood sanded 1/2 inch thick, of five layers, is used with the face plies running parallel to longitudinal direction, in accordance with the specification provided by the Douglas Fir Plywood Association.

When plywood is subject to a tensile or compressive force parallel to its length, only those plies having their grains running parallel to its length are considered as contributing to load carrying capacity. The cross plies are stressed across the grain and thus are incapable of contributing any significant amount to the strength of plywood in tension or compression. Similarly, in flexure, only the plies running parallel to the span are considered in computing the moment of inertia of the cross section (3, Section Two, p. 1).

Neglecting the moment of inertia of the face plies when plywood is used in flexure with its face grain perpendicular to the span,

$$I_1 = \frac{1}{12} \left[ \left( \frac{11}{32} \right)^3 - \left( \frac{3}{32} \right)^3 \right] = 0.00333 \text{ in}^4$$

### B. Ordinary Beam Approach

For a total load of 792 pounds uniformly distributed over the plate CC',

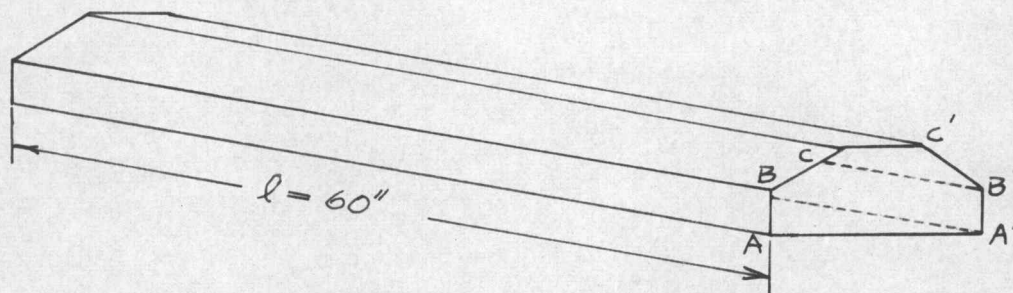


Fig. 2 - Hipped plate Roof

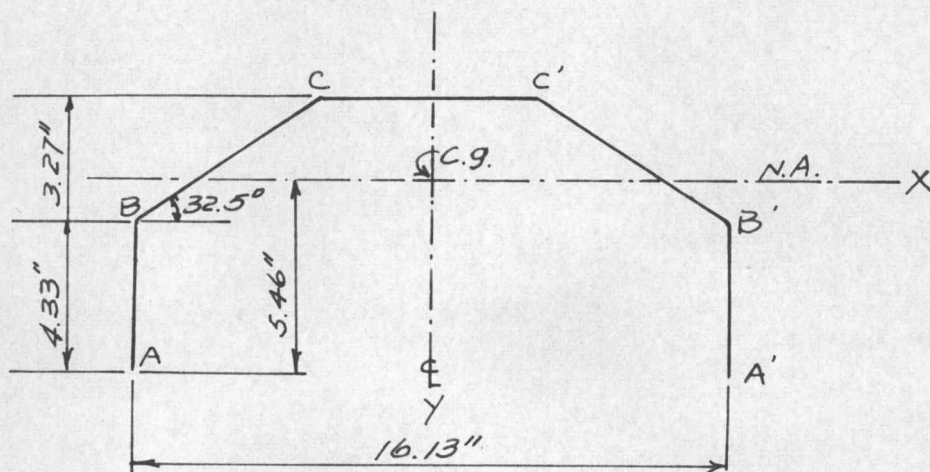


Fig. 3 - Cross Section of Roof

$$I_{xx} = 37.32 \text{ in}^4$$

$$w, \text{ load/linear inch,} = 792/60 = 13.2 \text{ lbs/in}$$

$$\text{Maximum Moment} = \frac{wl^2}{8} = 5,930 \text{ lb-in}$$

Maximum Unit Stresses on Top Fiber,

$$\begin{aligned} &= - \frac{M_{\max} C}{I_{xx}} \\ &= - \frac{(5,930)(2.14)}{37.32} \\ &= - 340 \text{ psi} \end{aligned}$$

Maximum Unit Stresses on Bottom Fiber,

$$\begin{aligned} &= + \frac{(340)(5.46)}{2.14} \\ &= + 867 \text{ psi} \end{aligned}$$

Maximum Unit Stresses on the Top of the Edge Beam,

$$\begin{aligned} &= + \frac{(340)(1.13)}{2.14} \\ &= + 180 \text{ psi} \end{aligned}$$

There is no doubt that the applicability of ordinary beam approach for this hipped plate model is doubtful

because of the lack of complete rigidity of the joints, the bending (out of plane) of the edge beam, and in part the nonhomogeneity of the material.

C. Theoretical Analysis by Hipped Plate Theory

The brief design procedure for the symmetrical and simply supported hipped plate structure is as follows:

1. Assume that the joints of the plates are held by the unyielding elastic supports.
2. Cut one inch center strip transversely and treat this strip to act as a continuous one-way slab.
3. Compute the transverse moments at the joints which cause the joint loads. Total joint loads are found by adding the joint loads by transverse moment to the joint loads by the applied load on the plate.
4. Resolve these joints loads into two components parallel to the adjacent plates to get the plate loads. Compute the longitudinal plate stresses.
5. Correct the longitudinal plate stresses found in Step 4. This correction is provided by the

requirement for compatible stresses in two adjacent plates at their common edge.

6. Assuming that the outer edge plate can rotate freely, compute the transverse moments due to the plate rotations. Repeat the Steps 3 to 5 to get the longitudinal stresses.
7. Final longitudinal plate stresses are found by adding the longitudinal plate stresses by Step 5 and by Step 6.
8. Transverse moments are superimposed similarly as in Step 7.
9. Compute the deflections perpendicular to the plates.
10. Using deflections found in Step 9. Compute the longitudinal stresses (out of plane) at edges.

The hipped plate roof shown in Fig. 4 will be analyzed. A total load of 792 pounds is placed on the top plate only. Neglecting the dead load,  $w' = 792 / (60)(6) = 2.2 \text{ lb/in}^2$  is assumed as live load.



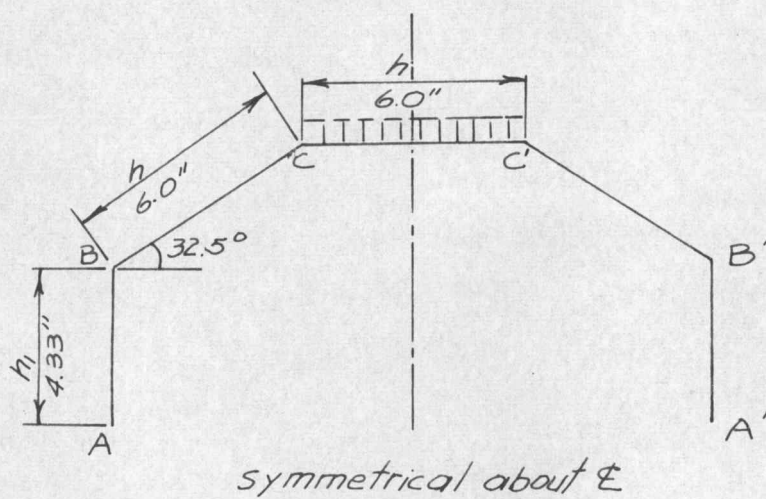
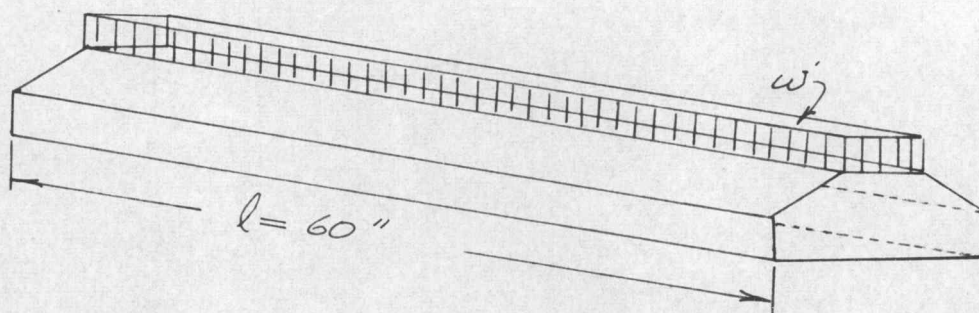


Fig. 4 — Proposed Plywood Model

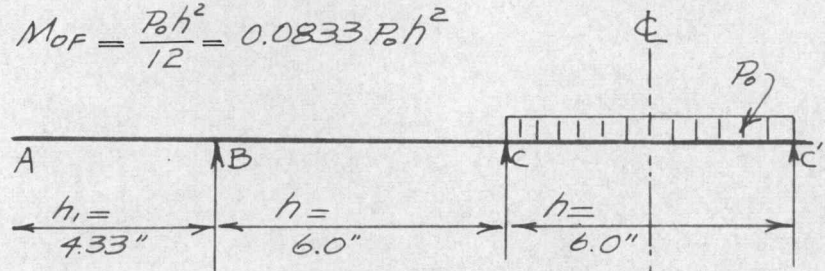
D. Case I

Assuming the Unyielding Elastic Support of  
A Continuous One-way Slab (transverse)

For 1 inch center strip.

$$p_0, \text{ line load, } \frac{16}{12} \text{ in/in} = \frac{4w'}{\pi} \sin \frac{\pi x}{l} = \frac{4w'}{\pi}$$

$$M_{OF} = \frac{P_0 h^2}{12} = 0.0833 P_0 h^2$$



Relative  $R$   
 Adjusted  $R$   
 FEM

(1)  
 (3/4)

(1)  
 (1/2)

$$-0.0833 P_0 h^2 + 0.0833 P_0 h^2$$

$$+0.05 P_0 h^2 + 0.0333 P_0 h^2 - 0.0333 P_0 h^2$$

$$+0.05 P_0 h^2 - 0.05 P_0 h^2$$

Moment

Shear

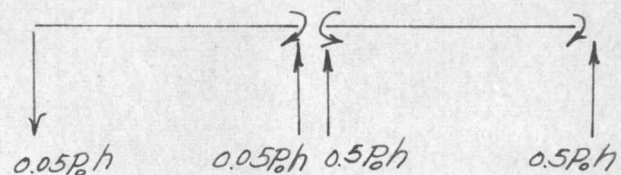


Fig. 5 — Transverse Moment and Shear



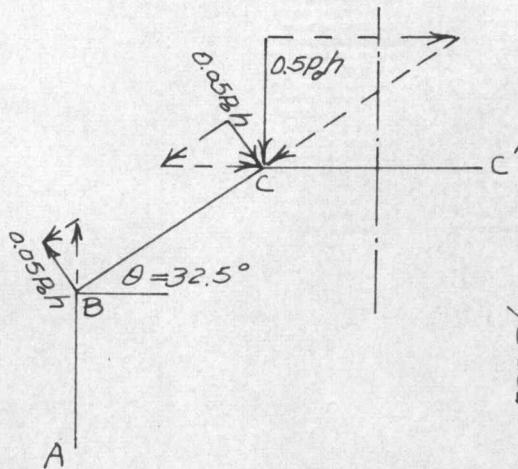
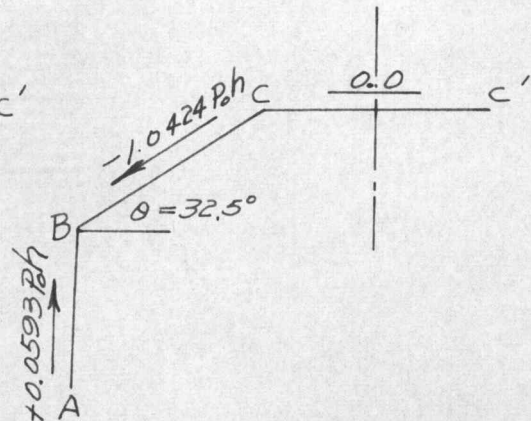


Fig. 6 - Joint Loads and Its Components

Fig. 7 -  $Q_0$ , Plate Loads

Sign convention: Clockwise direction is positive, and tension is positive.

Since plate loads vary as a half sine wave longitudinally, namely  $Q = Q_0 \sin \frac{\pi x}{l}$ ,

the longitudinal moment,  $M$

$$= \frac{l^2}{\pi^2} Q_0 \sin \frac{\pi x}{l}$$

$$M_0 = \frac{l^2}{\pi^2} Q_0$$

$$M_{0AB} = \left( \frac{l^2}{\pi^2} \right) (-1.0424 P_0 h)$$

$$M_{0BC} = \left( \frac{l^2}{\pi^2} \right) (0.0593 P_0 h)$$

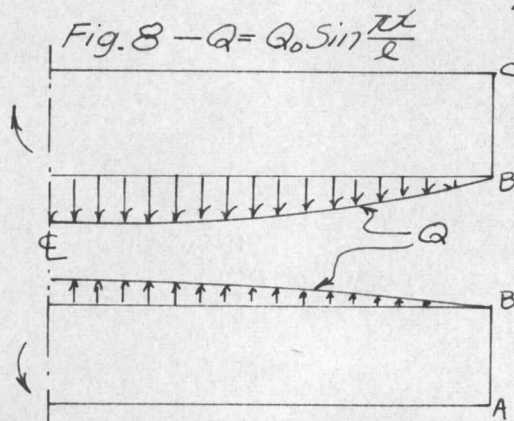


Plate	$M_0$ , lb-in	Section Modulus, $Z$ , in <sup>3</sup>	Stresses, psi
AB	$(l/\pi)^2 (0.0593 P_0 h)$	$Z = 0.782$	$\pm f_1 = 465$
BC	$(l/\pi)^2 (-1.0424 P_0 h)$	$Z = 1.5$	$\pm f_2 = 9,260$
CC'	0	$Z = 1.5$	0

TABLE 1 - Computation of the Longitudinal Stresses

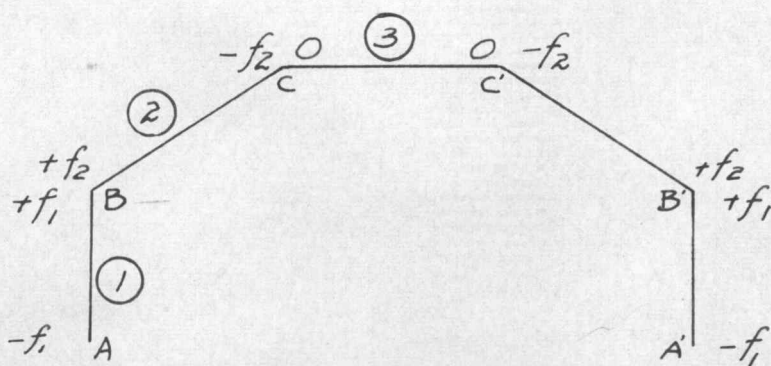


Fig. 9 — Longitudinal Edge Stresses

Due to unbalanced longitudinal edge stresses, longitudinal sliding would occur at the joint of two adjacent plates. In order to prevent this sliding, two equal opposite forces are introduced to produce the longitudinal edge shears.

Directions of these opposite forces are determined according to the magnitude of each edge stress.

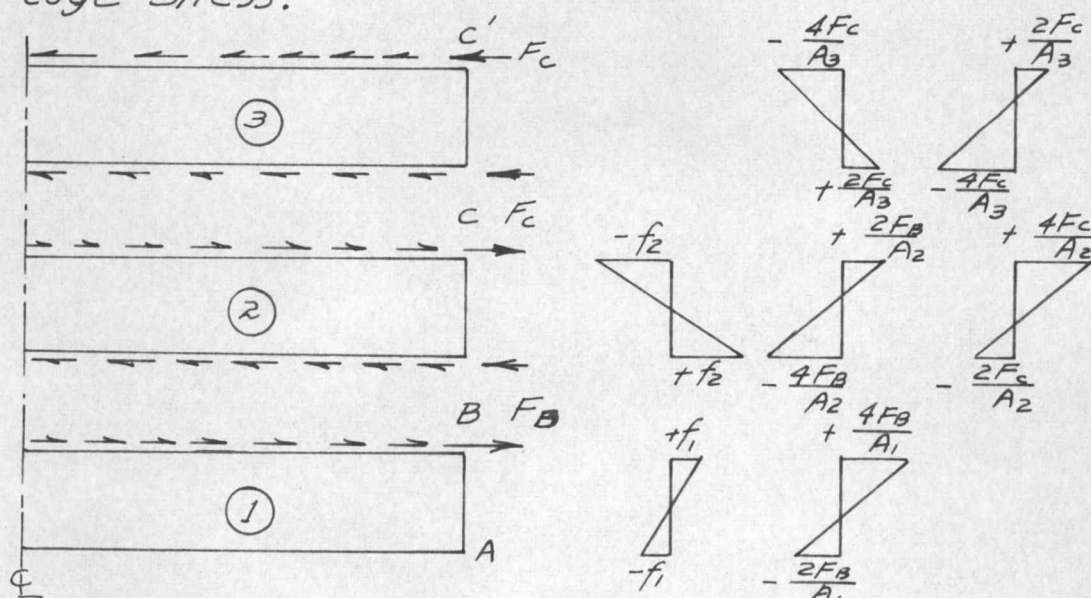


Fig. 10 — Edge Stresses by Edge Forces

Since the edge stresses are equal along the joint of two adjoining plates,

at joint C

$$\frac{4F_c}{A_2} + \frac{2F_B}{A_2} - f_2 = -\frac{2F_c}{A_3} \quad (2)$$

at joint B

$$-\frac{2F_c}{A_2} - \frac{4F_B}{A_2} + f_2 = \frac{4F_B}{A_1} + f_1 \quad (3)$$

Solving Eqs. (2) and (3),

$$F_c = \frac{t(26f_1 + 98f_2)}{115.3}$$

$$F_B = \frac{tf_2(3)(115.3) - 3t(26f_1 + 98f_2)}{115.3}$$

Substituting the values of  $F_c$  and  $F_B$  into Eqs. (2) and (3), the corrected longitudinal stresses are found.

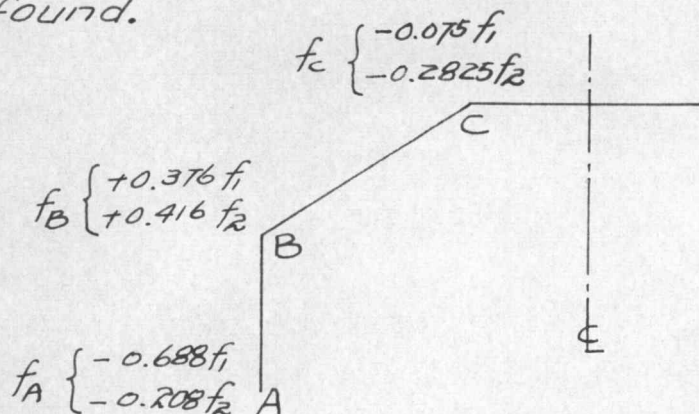


Fig. 11—Corrected Longitudinal Stresses



E. Case II

Considering plate Rotations

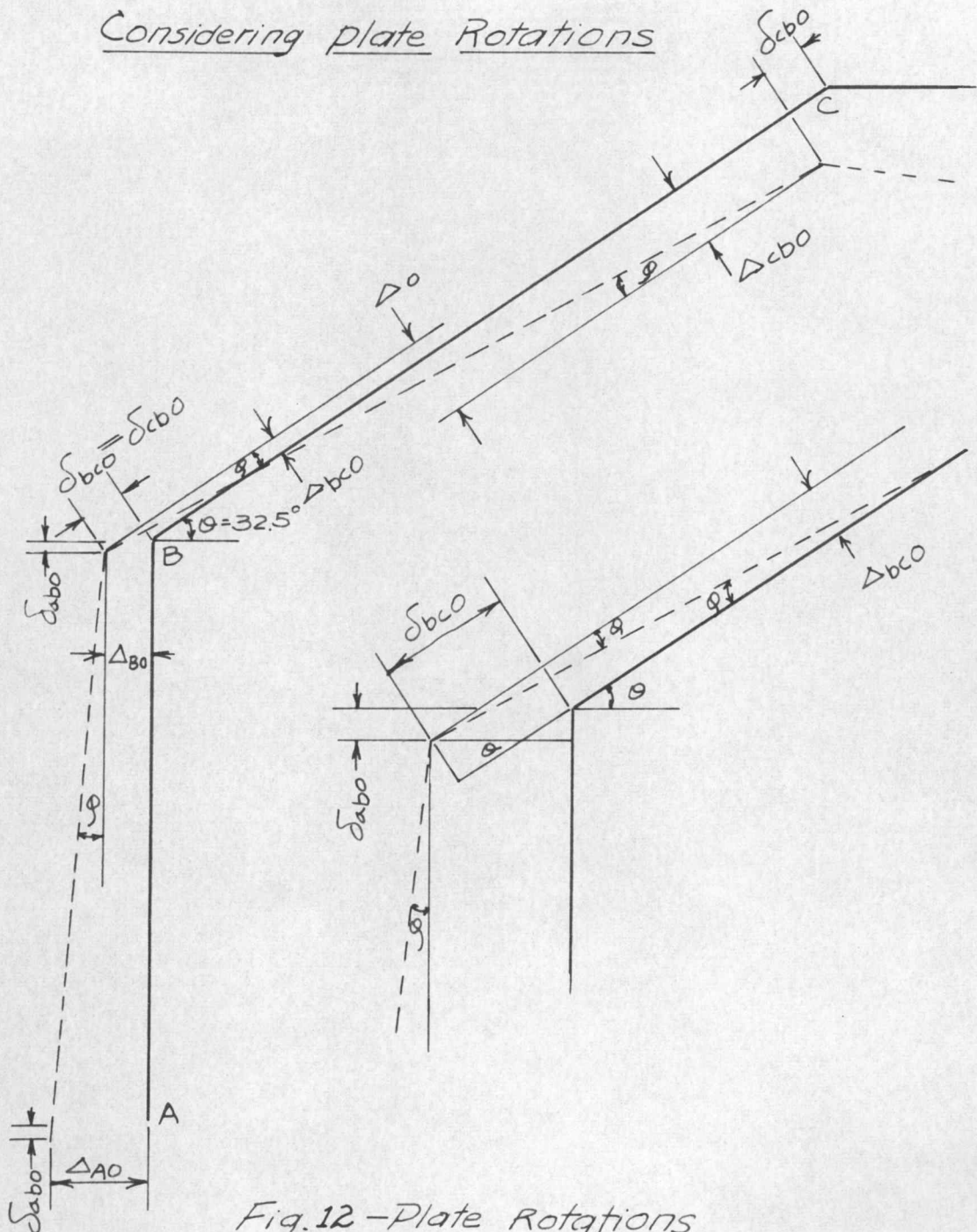


Fig.12 - Plate Rotations

Assuming the outer plate, AB, to rotate freely,

$$\text{F.E.M. at C, } M_{CF} = \frac{3EI}{h}(0 - \phi) = -\frac{3EI\Delta_0}{h^2}$$

$$= -3Y \quad \left( Y = \frac{EI\Delta_0}{h^2} \right)$$

Adjusted R

F.E.M.

Moment

Shear

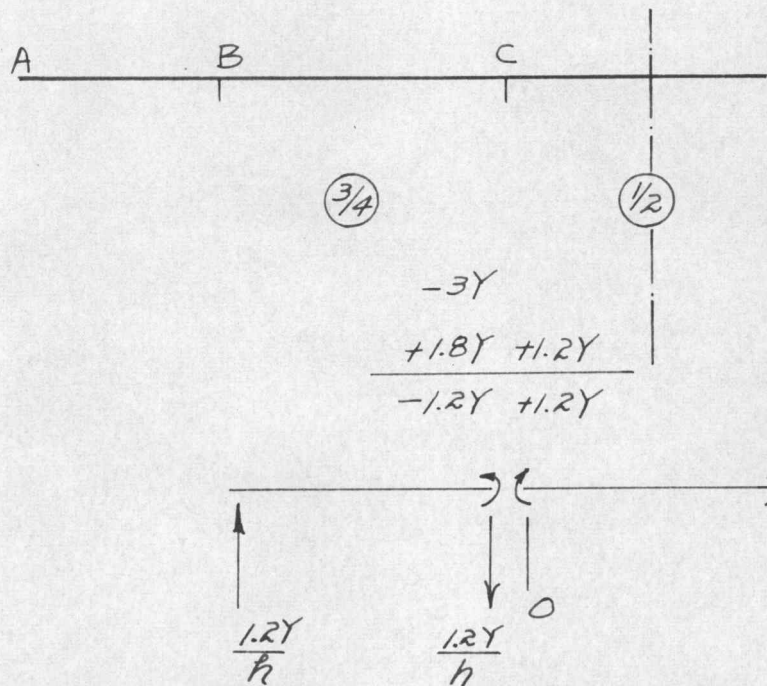


Fig. 13 — Transverse Moment and Shear

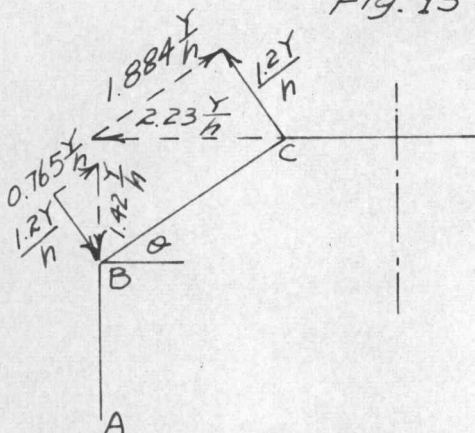


Fig. 14 — Joints Loads and Its Components

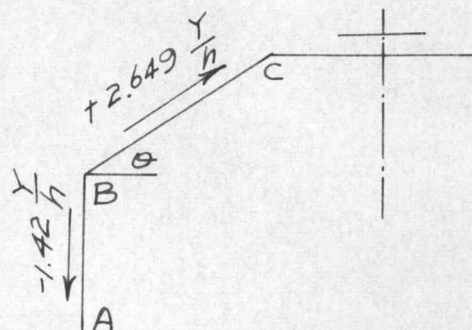


Fig. 15 —  $Q_0$ , Plate Loads

If the relative deflection " $\Delta$ " is assumed to vary as a sine wave, plate loads also vary as a sine wave.

$$\text{Plate loads, } Q' = Q_0' \sin \frac{\pi x}{l}$$

$$M = -\frac{l^2 Q_0'}{\pi^2} \sin \frac{\pi x}{l}$$

$$\delta = \frac{l^4 Q_0'}{\pi^4 EI} \sin \frac{\pi x}{l}$$

$$M_0 = -\frac{l^2 Q_0'}{\pi^2} \sin \frac{\pi \frac{l}{2}}{l} = -\frac{l^2 Q_0'}{\pi^2}$$

Plate	$M_0$ , lb-in	Section Modulus, $Z$ , in <sup>3</sup>	Stresses, psi
AB	$(Q/\pi)^2 (1.42 Y/h)$	$Z = 0.782$	$\pm f_1' = 110.7 Y$
BC	$(Q/\pi)^2 (-2.649 Y/h)$	$Z = 1.5$	$\pm f_2' = 107.5 Y$
CC'		$Z = 1.5$	0

TABLE 2 - Computation of the Longitudinal Stresses

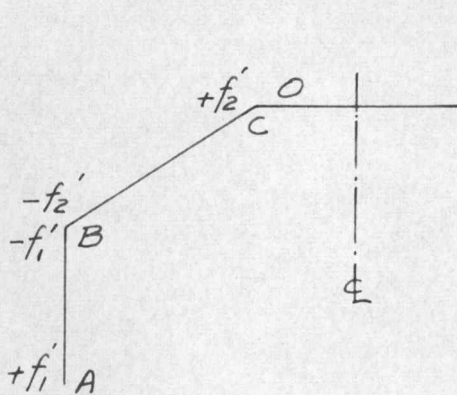


Fig. 16 - Longitudinal Stresses

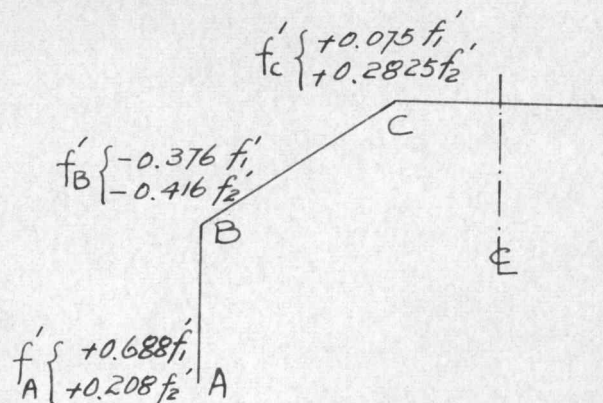


Fig. 17 - Corrected Longitudinal Stresses



### F. Deflection Formulas

For plate AB,

$$\delta_i = \frac{l^4 Q_0}{\pi^2 EI} \sin \frac{\pi x}{l}$$

$$\delta_{i0} = \frac{l^2 M_0}{\pi^2 EI} = \left( \frac{l^2}{\pi^2 EI} \right) \left( \frac{fI}{C} \right)$$

$$= \left| \frac{l^2 (f_A - f_B)}{\pi^2 E h_i} \right| \uparrow$$

$$\delta'_{i0} = \left| \frac{l^2 (f'_A - f'_B)}{\pi^2 E h_i} \right| \downarrow$$

The values of the deflection " $\delta$ " are absolute, and arrows designate the direction or sign.

These arrows are replaced by positive or negative signs according to the assumed direction of the deflected structure.

$$\delta_{abo} = -|\delta_{i0}| + |\delta'_{i0}|$$

$$= \frac{l^2}{\pi^2 E h_i} \left[ |f'_A - f'_B| - |f_A - f_B| \right] \quad (4)$$

Similarly for plate BC,

$$\delta_{bco} = \frac{l^2}{\pi^2 E h} \left[ |f_B - f_C| - |f'_B - f'_C| \right] \quad (5)$$

$$\Delta_{bco} = \tan \theta \left( \sigma_{bco} - \frac{\sigma_{abo}}{\sin \theta} \right) \quad (6)$$

$$\Delta_{cbo} = \sigma_{cbo} \cot \theta \quad (7)$$

$$\begin{aligned} \Delta_o &= \Delta_{bco} + \Delta_{cbo} \\ &= \sigma_{cbo} (\tan \theta - \cot \theta) - \frac{\sigma_{abo}}{\cos \theta} \\ &= \frac{l^2}{\lambda^2 E} \left[ (|f_B - f_c| - |f'_B - f'_c|) \left( \frac{\tan \theta + \cot \theta}{h} \right) \right. \\ &\quad \left. - \left( \frac{1}{\cos \theta h_1} \right) (|f'_A - f'_B| - |f_A - f_B|) \right] \quad (8) \end{aligned}$$

G. Computation of the Final Longitudinal Stresses

For case II,

$$\left. \begin{aligned} f'_1 &= \pm \left( \frac{l^2}{\lambda^2} \right) (1.42 \frac{Y}{h} \left( \frac{1}{Z_{AB}} \right)) \\ f'_2 &= \pm \left( \frac{l^2}{\lambda^2} \right) (2.649 \frac{Y}{h} \left( \frac{1}{Z_{BC}} \right)) \end{aligned} \right\} \quad (9)$$



Using Eq. (8),

$$\begin{aligned}\Delta_o &= \frac{l^2}{\pi^2 E} \left[ \left( \frac{\tan 32.5^\circ + \cot 32.5^\circ}{6} \right) (0.451 f_1 + 0.6985 f_2 \right. \\ &\quad \left. - 0.451 f_1' - 0.6985 f_2') - \left( \frac{1}{\cos 32.5^\circ} \right) \left( \frac{1}{4.33} \right) \right. \\ &\quad \left. (1.064 f_1' + 0.624 f_2' - 1.064 f_1 - 0.624 f_2) \right] \\ &= \frac{l^2}{\pi^2 E} (2.035 - 0.458 f_1' - 0.428 f_2') \quad (10)\end{aligned}$$

Substituting  $Y = \frac{EI \Delta_o}{h}$  into Eq. (9),

$$\begin{aligned}\Delta_o &= \frac{(\pi^2 f_1' h^3 E_{AB})}{(l^2)(1.42 EI_L)} \\ \frac{f_1'}{f_2'} &= \frac{1.42 E_{BC}}{2.649 E_{AB}} = 1.03 \quad (11)\end{aligned}$$

Solving Eqs. (10) and (11),

$$f_1' = 1,820 \text{ psi}, \quad f_2' = 1,770 \text{ psi}$$

The final longitudinal stresses are found by superposing the corrected longitudinal stresses of two cases.

However,  $|f_A - f_B| = 3,118 \text{ psi}$  is larger than  $|f_A' - f_B'| = 3,059 \text{ psi}$ . Therefore,  $\sigma_{abo}$  becomes negative sign which does not agree with the assumed deflected structure in Fig. 12.



From the geometric relations in Fig. 18,

$$\Delta_{bco} = \left( \delta_{bco} + \frac{\delta_{abo}}{\cos 57.5^\circ} \right) \cot 57.5^\circ \quad (6)'$$

$$\Delta_{cbo} = \delta_{bco} \cot \theta \quad (7)$$

$$\delta_{abo} = \frac{l^2}{\pi^2 E h_1} \left[ 3,151 - 0.624 f_2' - 1.064 f_1' \right]$$

$$= \frac{l^2}{\pi^2 E h_1} \left[ 3,151 - 185.1 Y \right]$$

$$\delta_{bco} = \frac{l^2}{\pi^2 E h} \left[ 3,184 - 125.0 Y \right]$$

$$\Delta_o = \Delta_{bco} + \Delta_{cbo}$$

$$= \frac{l^2}{\pi^2 E} \left[ \frac{1.186}{h_1} (3,151 - 185.1 Y) \right. \\ \left. + \frac{2.2071}{h} (3,184 - 125.0 Y) \right]$$

Since  $Y$  equal to  $\frac{E I_L \Delta_o}{h^2} = 148 \Delta_o$ ,

$$\Delta_o = \frac{l^2}{\pi^2 E} \left[ 2,035 - 14,330 \Delta_o \right]$$

$$= 0.1087 \text{ in.}$$

$$Y = 16.1 \text{ in.}$$



The final longitudinal stresses are given Table 3.

Edge	Case I, $f$ Psi	Case II, $f'$ Psi	Final Stresses (Case I + Case II) Psi
C	-1,237	+621	-618
B	+1,945	-1,388	+557
A	-1,206	+1,581	+375

TABLE 3—Final Longitudinal Stresses

#### H. Deflection Computations

$$\begin{aligned}\delta_{abo} &= \frac{l^2}{\pi^2 E h_1} [ |f_A - f_B| - |f'_A - f'_B| ] \\ &= \frac{(60)^2}{(\pi^2)(1.6 \times 10^6)(4.33)} [182] = 0.0096 \text{ in. } \uparrow\end{aligned}$$

$$\delta_{bco} = \frac{l^2}{\pi^2 E h_1} [ |f_B - f_C| - |f'_B - f'_C| ] = 0.0447 \text{ in. } \downarrow$$

$$\Delta_{co} = \frac{\delta_{bco}}{\sin 32.5^\circ} = 0.0832 \text{ in.}$$

$$\Delta_{Bo} = \frac{\delta_{abo}}{(\sin 32.5^\circ)(\cos 32.5^\circ)} = 0.0212 \text{ in.}$$

$$\Delta_{Ao} = \Delta_{Bo} + h_1 \phi = 0.0212 + (h_1) \left( \frac{\Delta_o}{h} \right) = 0.0999 \text{ in.}$$

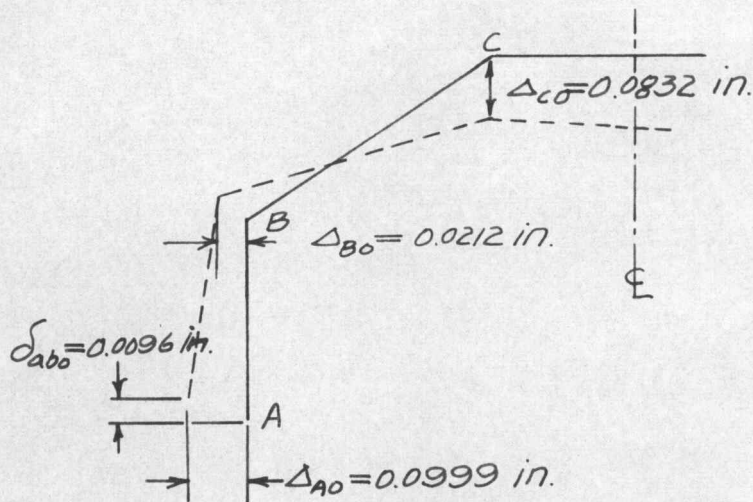


Fig. 19 - Displacements at Middle Section

I. Longitudinal Stresses (out of plane) at the Edges of Plate

Assuming that the deflections perpendicular to the plates vary as a half sine wave longitudinally, unequal longitudinal (out of plane) stresses at each joint will be found.

At joint C,

$$\Delta_{cb} = \Delta_{cbo} \sin \frac{\pi x}{l} = (\delta_{bco} \cot 32.5^\circ) \sin \frac{\pi x}{l}$$

$$M_o = EI \Delta_{cbo} \frac{\pi^2}{l}$$

$$f_{CB} = \frac{M_o c}{I} = \frac{(1.6 \times 10^6)(0.07)(\pi^2)}{(4 \times 60)^2} = 76.5 \text{ psi}$$



$$f_{cc'} = \frac{(f_{cB})(\Delta_{co})}{(\Delta_{cbo})} = \frac{(76.5)(0.0832)}{(0.07)} = 91.0 \text{ psi}$$

At joint B,

$$f_{Bc} = (f_{cB}) \left( \frac{\Delta_{bco}}{\Delta_{cbo}} \right) = (f_{cB}) \left( \frac{\Delta_o - \Delta_{cbo}}{\Delta_{cbo}} \right) = 42.3 \text{ psi}$$

$$f_{BA} = (f_{cB}) \left( \frac{\Delta_{Bo}}{\Delta_{cbo}} \right) = (76.5) \left( \frac{0.0212}{0.07} \right) = 23.2 \text{ psi}$$

At edge A,

$$f_{AB} = (f_{cB}) \left( \frac{\Delta_{Ao}}{\Delta_{cbo}} \right) = (76.5) \left( \frac{0.0999}{0.07} \right) = 108.0 \text{ psi}$$

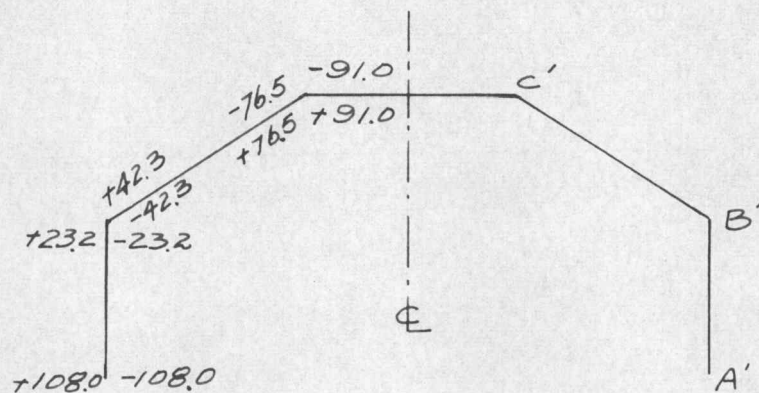


Fig. 20 - Longitudinal (out of plane) Stresses, psi

#### J. Transverse Moments

The bending moments in a transverse 1.0 inch wide strip are found by superposing the moments in two cases. In case I, in which

external loads and unyielding supports considered, the connecting moment at joint C is  $M_{co} = 0.05 P_o h^2 = 5.04 \text{ lb-in.}$  The maximum moment with uniform load in span CC' is  $\frac{P_o h^2}{8} = 12.6 \text{ lb-in.}$  In case II, the moment due to the plate rotations is  $M'_{co} = 1.2 Y = 19.3 \text{ lb-in.}$

Joint	Case I Lb-In.	Case II Lb-In.	Final Moments (Case I + Case II) Lb-In.
C	5.04	-19.3	-14.26
B	0	0	0

TABLE 4 — Transverse Moments in 1.0 inch Strip

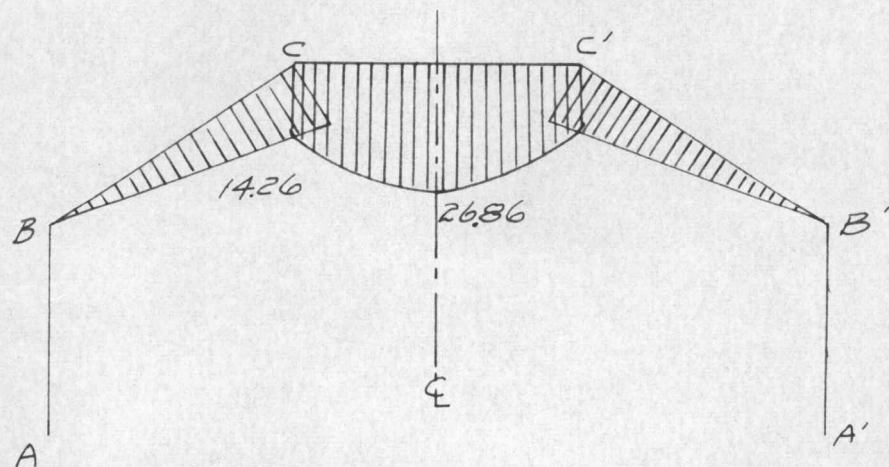


Fig. 21 — Transverse Moments, Lb-In.

## PART IV

EXPERIMENTAL INVESTIGATION  
OF A PLYWOOD MODELA. Model Description

The material descriptions and strength properties are mentioned in p. 9. For loading six copper plates weighing 5.5 pounds each were placed, as one layer, over the top horizontal plate. These layers were stacked until the desired weight was loaded as shown in Fig. 22.

The edges of plates at the joints are not rectangular shape and dimensions are measured along the centerline of the plate cross section as shown in Fig. 24. The plates were glued at the joints and nailed at six inch intervals. The bottoms of the end diaphragms were also glued and nailed in order to prevent "tip-in" action which would cause the end diaphragms to bend inward as the load is applied. Ames dial gages were installed at edges A and C to determine vertical deflections; at edges A, B, and quarter point of the bottom of the edge beam for horizontal deflections; and at the top of end diaphragm for horizontal "tip-in" action. SR-4 strain gages

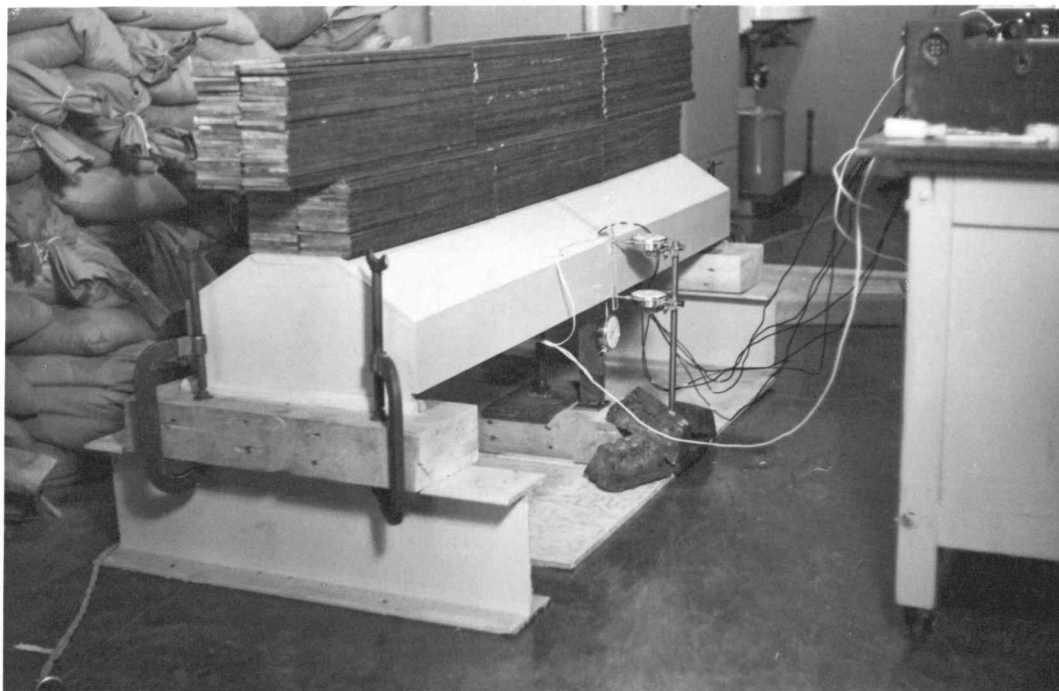


Figure 22

Loaded Model with Gages



were placed at the joints, free edge, and center of the plate CC' for longitudinal strains; a transverse SR-4 strain gage was placed at the bottom of the joint C as shown in Fig. 24. In Fig. 23, 1, 2, 3, 4, 5, 6, and 7 are the longitudinal strain gages and 8 the transverse gage. The picture in Fig. 22 does not show the Ames dial gages at the top of the end diaphragm and quarter point of the edge beam, nor does it show strain gages 2 and 6. These strain gages were wired independently to give the strain readings at outer and inner surfaces of the plate separately.

#### B. Experimental Procedure

The first experiment was made by increasing loads 100 pounds at a time up to a total of 800 pounds. Gage readings were recorded at 100 pound increments. In the second experiment, loads were applied at 200 pound increments up to 800 pounds, then at 50 pound increments up to a total of 1,950 pounds. At each loading interval the gage readings were taken. The third experiment with the loading increments of 200 pounds were run increasing the load to 800 pounds. From 800 pounds up to a total of

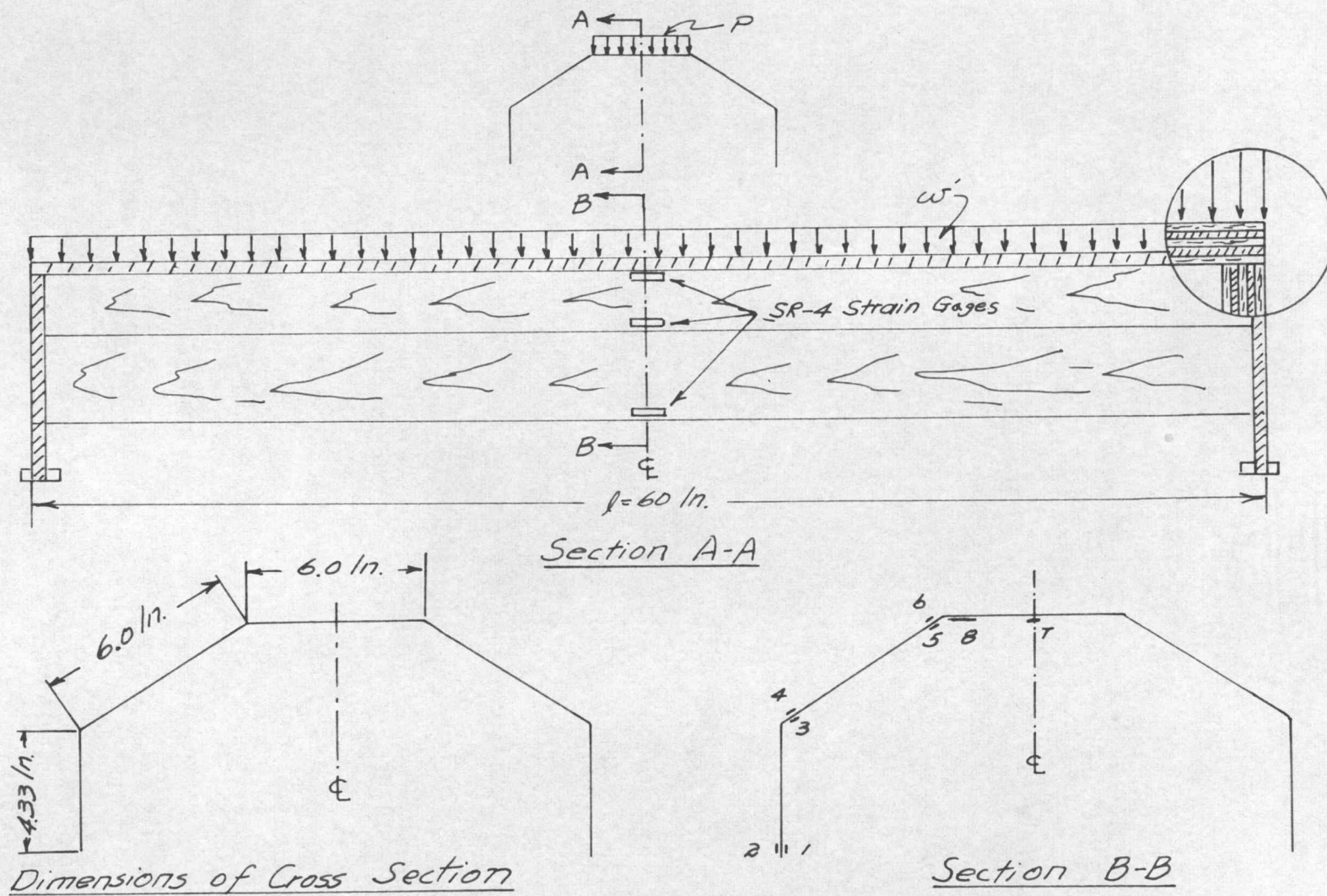


Fig. 23 - Plywood Model

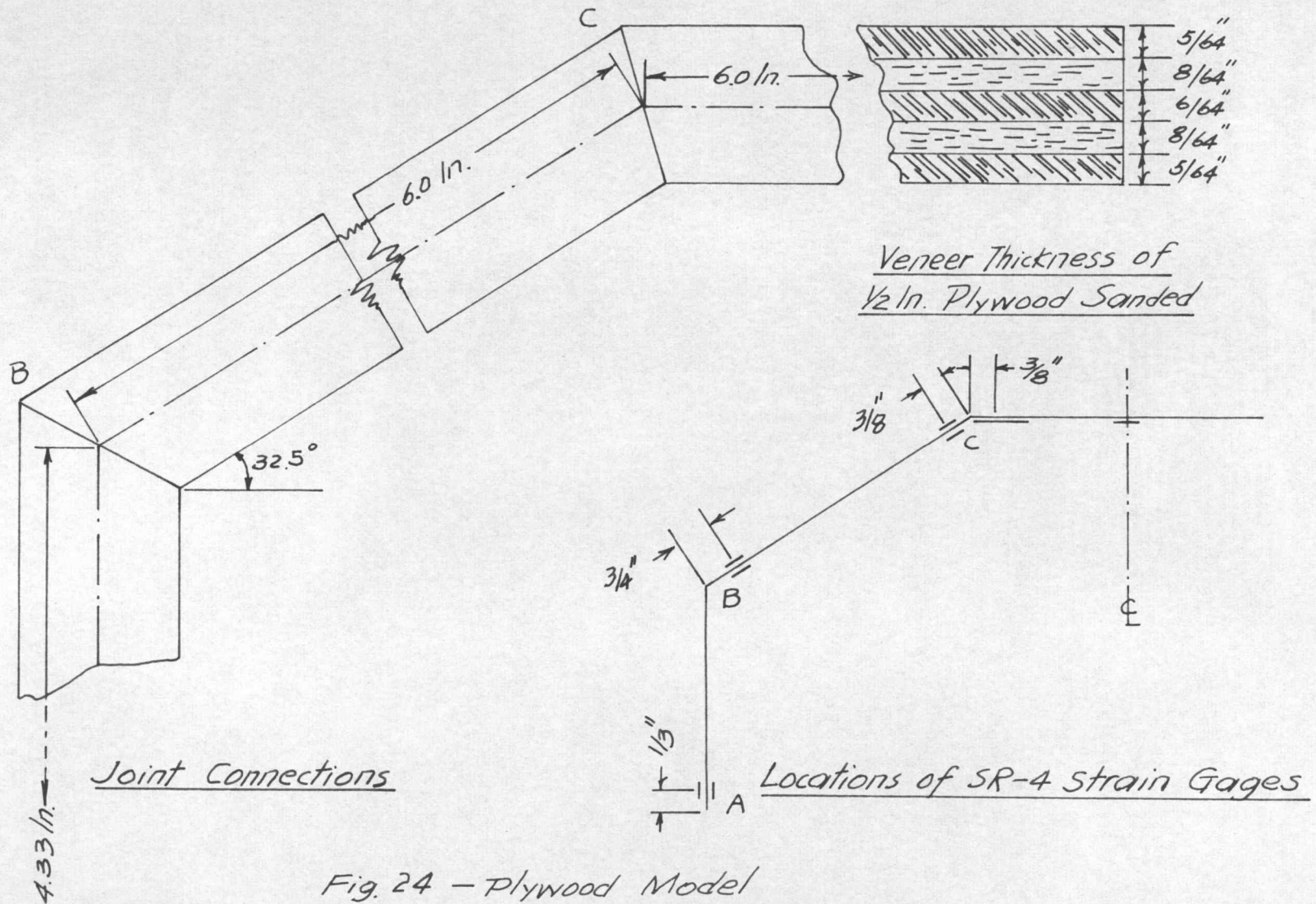


Fig. 24 - Plywood Model

1,600 pounds, the load was added at one time and the gage readings were taken only once.

### C. Results

#### 1. Observations of gage reading

Immediately after each loading, it was observed that both strain and Ames dial gage readings were creeping rapidly for about one and one-half minutes, then steadying. It is believed that this creep is caused by the fact that certain time is required for plastic flow and joint translations to occur. The gage readings in the experiment were taken after an elapsed time of approximately two minutes.

#### 2. Strain Measurements

The longitudinal and transverse strains at the middle section of the model are shown in Fig. 26. For the plate CC', a top surface strain gage was not installed because of the difficulty caused by the loading on the top of the plate. Longitudinal plate stresses were found from the average strain values of the inner and outer surface of the plate. A transverse strain of 350 micro-in./in. was measured



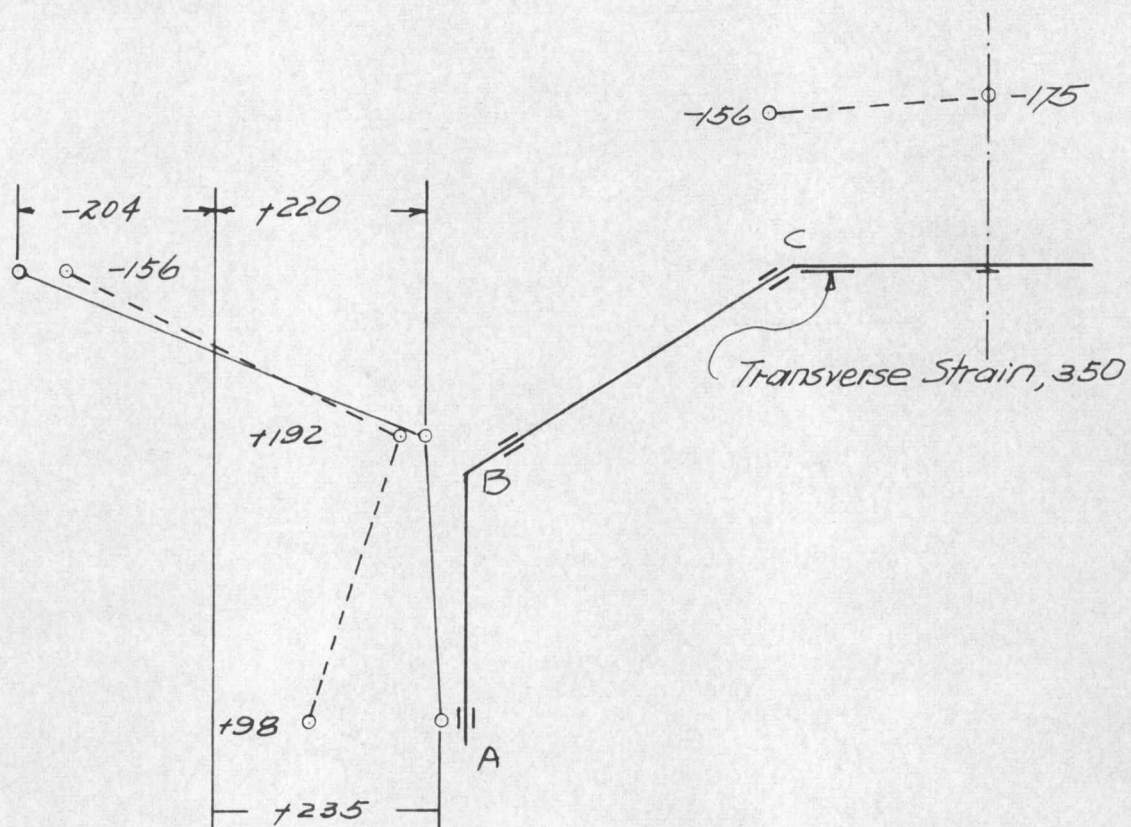


Fig. 25 - Maximum Longitudinal Strain,  $\frac{\text{Micro-Inch}}{\text{Inch}}$

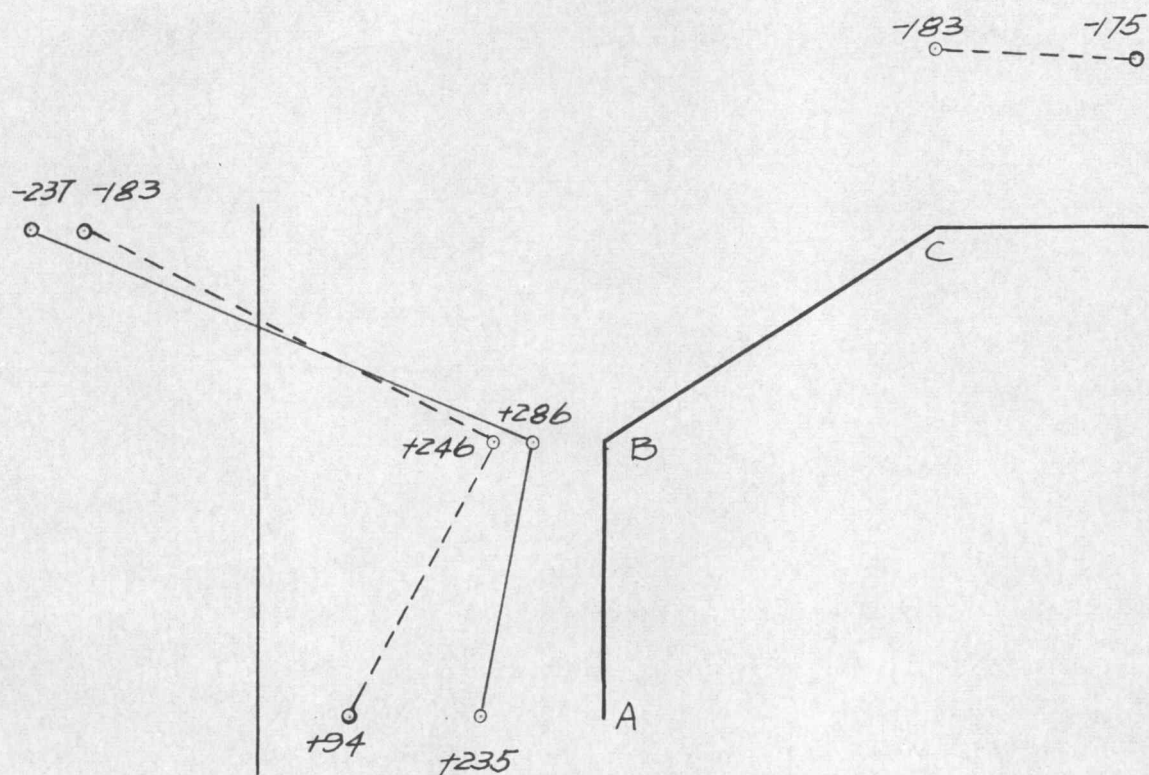


Fig.26 — Maximum Longitudinal Strain  
at the Joint ,  $\frac{\text{Micro-Inch}}{\text{Inch}}$

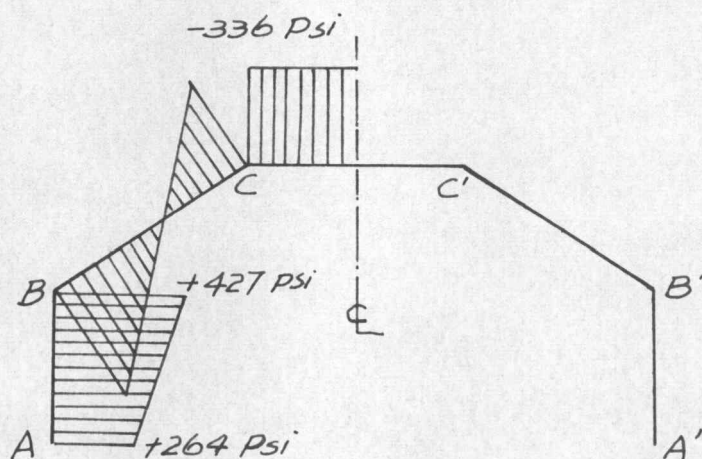
Note :

Dotted lines are for inner surface  
Solid lines are for outer surface

at joint C. Transverse moment in center strip of the model yields as the connecting moments at the joint

$$C, M_{co} = \epsilon E L / C = (350)(1.6)(0.0033) / 0.25 = 7.42 \text{ lb.-in.}$$

As the transverse deflection in span CC' is negligible due to the short span length compared to the longitudinal length of the model, this 7.42 lb.-in. is considered as an average moment.



*Fig. 27—Final Longitudinal Edge Stresses*

### 3. Displacement Measurements

The measured displacements of the edges are shown in Fig. 28. The values of  $\Delta_A$  and  $\Delta_{AO}$  verify the previous assumption that the deflections perpendicular to the plate vary as a half sine wave. The downward vertical displacement of the plate AB does



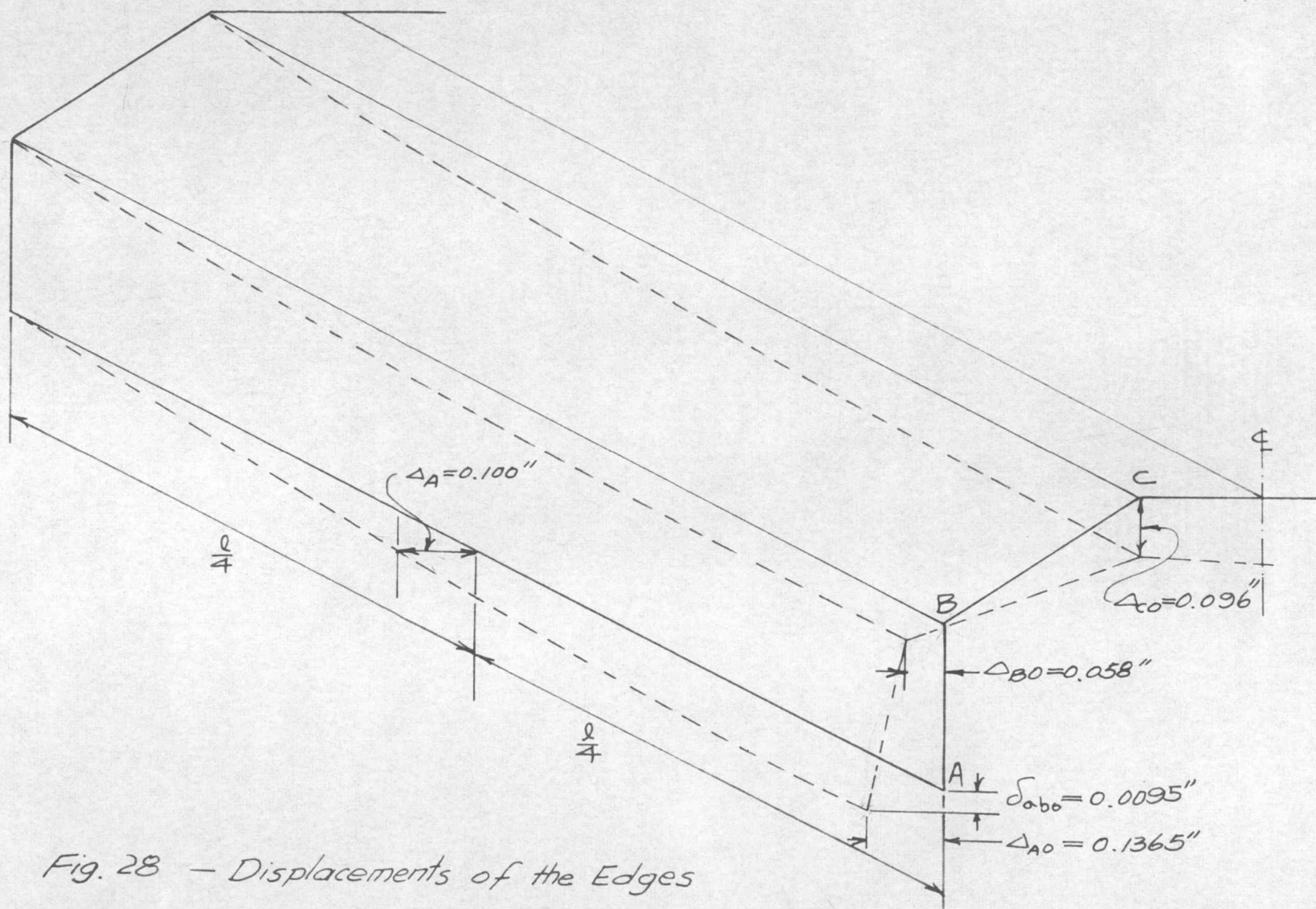


Fig. 28 — Displacements of the Edges



not agree with the theoretical displacement which is upward.

The possible causes of this disagreement which at present can be offered are, that the wooden block under the end diaphragm compressed as the load was applied and the glued joints were not rigid enough to translate the connecting moments fully. Since the longitudinal (out of plane) stresses are of the same order of magnitude as the deflections, it is not necessary to compute the longitudinal (out of plane) stresses for the comparison of theoretical and experimental values.

## PART V

DISCUSSIONS ON THE RESULTS OF THE  
EXPERIMENTAL AND THEORETICAL INVESTIGATION

In ordinary beam theory the longitudinal stresses are proportional to the distances from the centroid of the entire section, whereas they are not proportional in hipped plate theory.

(1) Edge	Stresses, Psi		
	(2) Ordinary Beam Theory	(3) Hipped Plate Theory	(4) Experimental Values
C	-340	-618	-336
B	+180	+557	+427
A	+867	+375	+264

Table 5 - Maximum Longitudinal Stresses

Edge	Plate	Types of Deflec- tions	Deflection by Hipped Plate Theory, Inch	Deflection by Test, Inch
A	AB	$\delta_{ab}$	0.0096	-0.0095
A	AB	$\Delta_{AO}$	0.0999	0.1365
B	BA	$\Delta_{BO}$	0.0212	0.058
C	CC'	$\Delta_{CO}$	0.0832	0.096

Table 6 - Deflections at the Middle Section

In other words, the hipped plate structures do not behave as a unit. Thus beam theory cannot be applied directly to the structures of nonhomogeneous material, not to the hipped plate structures. However, the beam theory gives the right sense of the signs of the stresses as shown in Table 5. The longitudinal stress values by hipped plate theory are much higher than the experimental values. Assuming that theoretical stresses in column (3) of Table 5 are the correct values, the errors of the longitudinal stresses at the edges A, B and C are:

$$\frac{375 - 264}{375} \times 100 = +30\%, \quad \frac{557 - 427}{557} \times 100 = 23\%$$

$$\text{and } \frac{618 - 336}{618} \times 100 = 46\%.$$

The possible causes for these errors are as follows:

- A. During the experiment, a slight twist of the model was observed.
- B. The material is nonhomogeneous, nonmonolithic, and nonisotropic.
- C. In theoretical analysis, strength property calculations of the plywood might not be exactly

correct, even though the model was designed in accordance with the specification by the Douglas Fir Plywood Association.

- D. The joints are probably not completely homologous.
- E. The glued connected joints were not stiff enough to be regarded as rigid joints.
- F. The cross sections of the plate at joints were not rectangular.
- G. The effect of nonuniform loading due to overhang of top bars and possible bridging.

The computed maximum transverse moment at joint C is 7.42 pound-inches by test and 14.26 pound-inches computed by the hipped plate theory. The displacements of the edges are listed in Table 6 for comparison. Since the transverse moments, longitudinal stresses, and deflections of the edges are of the same order of magnitude, the errors in transverse moments and deflections of the edges probably result from the above causes.

The deflections perpendicular to the plate are important in the hipped plate analysis. This type of deflection is the main cause of the different strains in the inner and outer surfaces of the plate. Since the



difference in longitudinal strains between the inner and outer surfaces is twice the longitudinal (out of plane) strains, from Fig. 26 the approximate longitudinal (out of plane) stresses can be computed. At the edges A, B, and C the longitudinal (out of plane) stresses are:

$$\left(\frac{235 - 94}{2}\right)(1.6) = 113, \quad \left(\frac{286 - 246}{2}\right)(1.6) = 32, \text{ and}$$

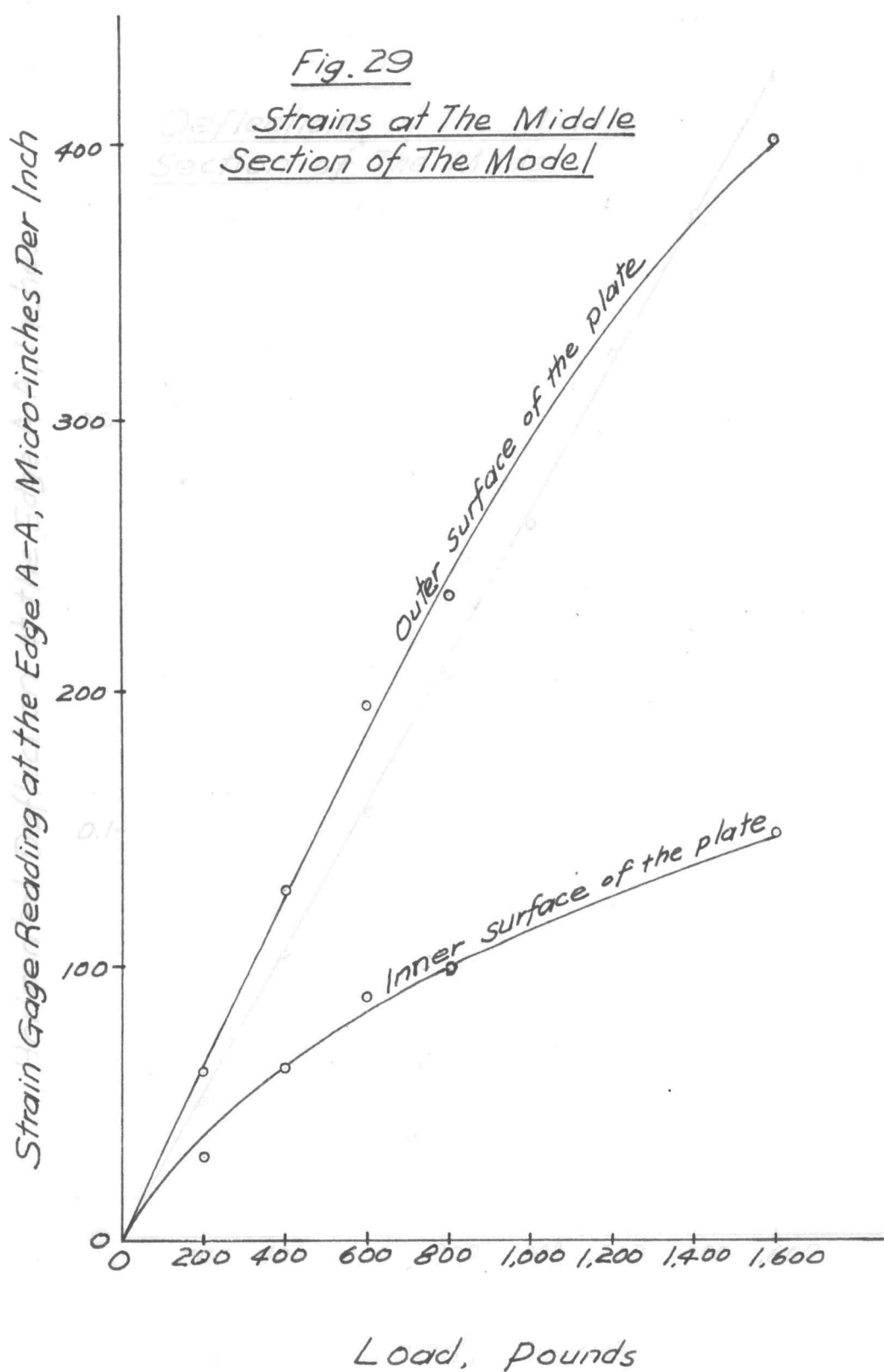
$\left(\frac{237 - 183}{2}\right)(1.6) = 43$  psi. These values agree quite well with the computed values of the theoretical analysis in Fig. 20 except at the edge C. By computation, assuming deflections as a sine wave, at the edge C the longitudinal (out of plane) stress is:

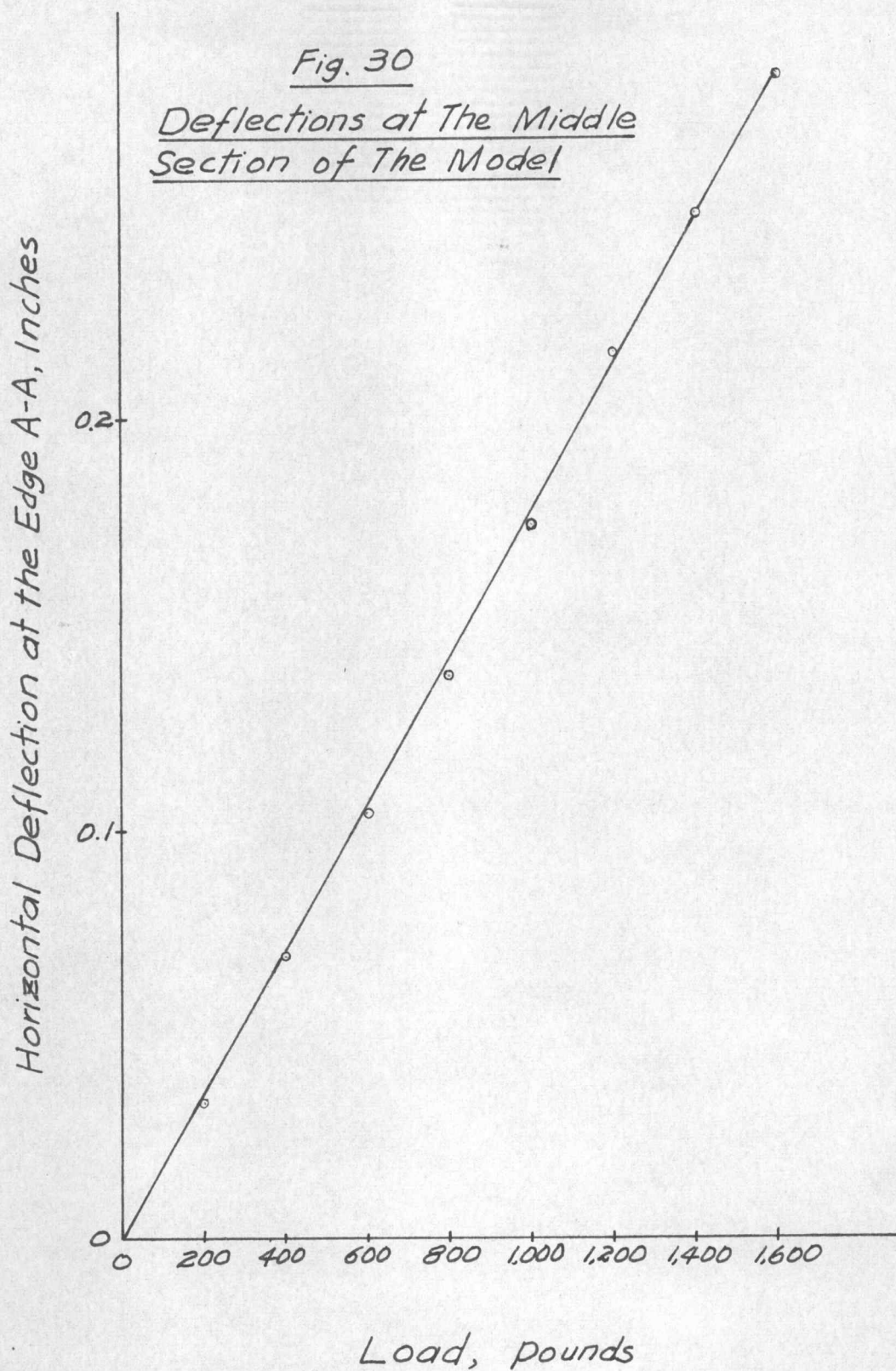
$$f_{cc'} = \frac{EI\Delta_{co}\pi^2c}{I\ell^2} = 104.0 \text{ psi}$$

The equation of the longitudinal (out of plane) stresses is expressed as:

$$f_{cc'} = \frac{EI\Delta\pi^2c}{I\ell^2} = R\Delta = R'(P_o, Q_o) = R''(P_o - Q_o) \quad (12)$$

The curves of the load versus strain at the joints are linear except at the edge A. This statement of the linear relationship between load and strain, agrees with the equation (12) by theory.





At the edge A, however, the load versus strain curve is not linear despite the linear relationship between load and deflection as shown in Fig. 29 and Fig. 30. In Fig. 29, up to about 1,000 pounds the load versus strain curve is linear. Beyond 1,000 pounds, in the writer's opinion, some unknown factors are complicating the method of analysis presented in this paper. Possibly this non-linearity beyond 1,000 pounds might be due to the same causes of the errors in the longitudinal stresses.



## PART VI

### CONCLUSIONS

If the properties of plywood were more uniform and more reliable data on plywood were available, a plywood model would be easy to build and satisfactory for studying the hipped plate structures. Rigid joint connections are also essential for perfect model construction.

Since  $\Delta_{A0} = \Delta_0 \sin \frac{\pi x}{l} = \phi(\delta, \Delta)$  is fairly well verified by the experiment, the assumption made in the theoretical analysis, in which the elastic line of the plate along the longitudinal edges is expressible as functions of  $\sin \frac{\pi x}{l}$ , is confirmed. As the load is increased the effects of the longitudinal (out of plane) stresses appear to be dominant.

It is believed that the proposed theory presented in this thesis describes more completely the structural action of hipped plate structures than most methods currently used in design. It is hoped that further testing may aid in verifying this theory and expanding upon obscure points.

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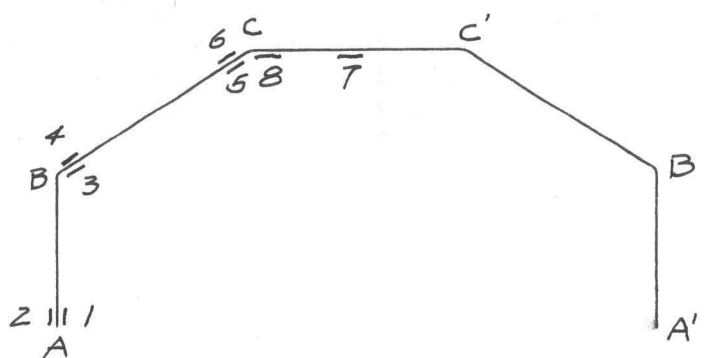
# A P P E N D I X A

## TEST DATA



**First Experiment**

Load Lbs.	Gage Readings, Micro-inches per Inch					
	Gage Number					
	1	3	4	5	7	8
0	5-1340	5- 830	5- 950	5-1140	5- 530	5- 740
100	1348	840	5- 978	1125	518	805
200	1360	870	5-1010	1113	500	860
300	1385	910	5-1030	1100	473	905
400	1385	940	5-1080	1075	460	950
500	1405	970	5-1092	1055	430	990
600	1405	997	5-1150	1030	420	1030
700	1420	1030	5-1160	1010	400	1060
800	1423	1043	5-1200	990	387	1090
Net Strain	+83	+213	+250	-150	-143	+350



Gages except No. 8 are longitudinal strain  
 Gage No. 8 is transverse strain

**Deflections at Middle Section  
of the Model**

Load Lbs.	Vertical Deflection Inches		Horizontal Deflection Inches	
	Joint		Joint	
	A	C	A	B
0	0.0005	0.017	0.0155	0.0025
100	0.002	0.0135	0.0185	0.0085
200	0.0025	0.014	0.0195	0.009
300	0.0015	0.015	0.0215	0.015
400	0.0010	0.016	0.0195	0.000
500	0.0010	0.016	0.013	0.005
600	0.0020	0.013	0.021	0.008
700	0.0010	0.012	0.017	0.0075
800				
Total	0.0115	0.1165	0.1455	0.0555

**Second Experiment**

Load Lbs.	Gage Readings, Micro-inches per Inch					
	Gage Number					
	1	3	4	5	7	8
0	1388	878	1000	1190	590	780
200	1412	930	1065	1158	546	875
400	1440	975	1120	1118	498	957
600	1456	1010	1160	1058	441	1022
800	1486	1060	1210	1035	420	1100
900	1500	1090	1235	1015	400	1140
1000	1505	1100	1260	1000	390	1160
1100	1500	1140	1295	970	375	1180
1200	1515	1158	1318	950	357	1210
1300	1523	1188	1323	922	335	1233
1400	1510	1208	1372	900	323	1243
1500	1520	1232	1380	870	305	1255
1600	1520	1242	1422	848	295	1275
1700	1522	1275	1430	820	278	1288
1733	1495	1285	1470	803	280	1273
1788	1498	1317	1478	782	270	1288
1843	1495	1320	1515	762	262	1295
1953	1500	1360	1527	730	237	1323
Net Strain	+112	+482	+527	-460	-353	+543

Load Lbs.	Vertical Deflection Inch		Horizontal Deflection Inch	
	A	C	A	B
0	0.003	0.0235	0.0325	0.012
200	0.002	0.025	0.0360	0.016
400	0.0025	0.0255	0.0350	0.016
600	0.0020	0.022	0.033	0.014
800	0.0005	0.0125	0.020	0.009
900	0.000	0.0115	0.017	0.007
1,000	0.001	0.015	0.019	0.008
1,100	0.000	0.013	0.023	0.011
1,200	0.001	0.011	0.017	0.007
1,300	0.000	0.015	0.023	0.010
1,400	0.0005	0.013	0.019	0.009
1,500	0.0003		0.0145	
1,600				



**Third Experiment**

Load Lbs.	Gage Readings, Micro-inches per Inch Gage Numbers							
	1	2	3	4	5	6	7	8
0	1392	690	860	992	1166	704	595	725
200	1422	750	902	1035	1120	641	525	782
400	1455	818	960	1100	1092	595	489	885
600	1480	885	1000	1142	1052	548	452	980
800	1490	925	1042	1192	1010	500	420	1022
	+98	+235	+182	+200	-156	-204	-175	+297
1600	1540	1092	1180	1350	850	350	292	1260
Net	+148	+402	+320	+358	-316	-354	-303	+535

Load Lbs.	Vertical Deflection Inch A	Horizontal Deflection Inch at the Top of End Diaphragm
0	0.005	0.00
200	0.0015	0.0025
400	0.0010	0.0005
600	0.0005	
	<u>0.008</u>	<u>0.0030</u>
800	0.006	
1,600		

Horizontal Deflections at Edge A-A

Load Lbs.	Center of the Edge Inches	Quarter Point of the Edge Inches
0	0.142	0.100
800		

A P P E N D I X B

NOTATIONS

## NOTATIONS

$A$	cross sectional area, ht
$A, B, CC', a, b, c$	subscripts used to denote the edges or joints
$E$	modulus of elasticity of plywood, $1.6 \times 10^6$ psi
$F$	edge forces
$f_1, f_2$	longitudinal stresses of the corresponding plates in case I
$f_A, f_B$	corrected longitudinal stresses
$f_{BC}, f_{CB}, f_{AB}, f_{BA}, f_{CC'}$	longitudinal (out of plane) stresses, acting on the edge (first subscript) of the corresponding plate
$h_1$	width of the plate AB
$h$	width of the plate BC or CC'
$I_{  }$	moment of inertia when the face grain is parallel to the span
$I_{\perp}$	moment of inertia when the face grain is perpendicular to the span
$k$	rigidity coefficient ( $k = EI/h$ )
$l$	span of the roof between end diaphragms
$o$	subscript used to denote the middle of a beam or plate
$p$	line load,



Q	plate load
t	thickness of plate
w	uniformly distributed load in beam theory analysis
w'	uniform distributed load in hipped plate theory analysis
Z	section modulus of a plate, considered to act as a beam
'	superscripts used to denote the values in case II, for example, $f_1'$ , $f_2'$ , $f_A'$ , $f_B'$ , $\delta_1'$ , etc.
1, 2, 3	subscripts used to denote plates: plate AB as 1, BC as 2, and CC' as 3
$\Delta$	deflection perpendicular to the plate and also $\Delta = \Delta_{bc} + \Delta_{cb}$
$\Delta_{bc}$ , $\Delta_{cb}$	deflection of the edge (first subscript) perpendicular to the corresponding plates
$\Delta_A$ , $\Delta_B$	deflection perpendicular to the plate AB, at the edges A and B
$\Delta_c$	deflection perpendicular to the plate CC' at the edge C
$\delta$	deflection in the plane of the plate
$\delta_{bc}$	plate deflection of the plate BC, $\delta_{bc} = \delta_{cb}$

$\delta_1$ 

plate deflection of the plate 1 or  
AB in case I

 $\delta_{ab}$ 

plate deflection of the plate AB,

$$\delta_{ab} = \delta_1 + \delta_1'$$

 $\varphi$ 

angle changes due to the plate  
rotations